Precision Higgs boson mass determination at photon-photon colliders

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(January 12, 1999)

Abstract

We demonstrate a measurement of the Higgs boson mass by the method of energy scanning at photon-photon colliders, using the high energy edge of the photon spectrum. With an integrated luminosity of 50 fb^{-1} it is possible to measure the standard model Higgs mass to within 110 MeV in photon-photon collisions for m_h = 100 GeV. As for the total width of the Higgs boson, the statistical error $\Delta\Gamma_h/\Gamma_{h\ \mathrm{SM}}$ = 0.06 is expected for m_h = 100 GeV, if both $\Gamma(h \to \gamma\gamma)$ and $\Gamma(h \to b\bar{b})$ are fixed at the predicted standard model value.

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One of the most important tasks of the current and the future collider experiments will be to detect and study Higgs boson(s). The very accuracy of the measurement of the Higgs boson mass will impact precision tests of loop corrections, both in the standard model (SM) and in the extended models such as the minimal supersymmetric model (MSSM) [1–3]. Deviations of the total widths of the Higgs bosons from SM predictions can be directly compared to predictions of alternative models such as the MSSM, the non-minimal supersymmetric standard model, or the general two-Higgs-doublet model [1–3]. The total widths for the SM Higgs boson h_{SM} and the three neutral Higgs bosons h^0, H^0, A^0 of the MSSM are shown in Fig. 1.

The interaction of high energy photons at a photon-photon collider [4–6] provides us with an unique opportunity to study Higgs boson, because the SM Higgs boson in s-channel resonance can be produced at photon-photon colliders [7–11]. In this letter we point out precision measurements of the mass (m_h) and total width (Γ_h) of the Higgs boson by the method of energy scanning, using the high energy edge of the photon spectrum.

The method of energy scanning at photon-photon colliders was first mentioned by V. Telnov [6]. The luminosity of the photon-photon collider has a very sharp edge at high energy, much narrower than the width of the luminosity peak. If the Higgs boson is a very narrow resonance, we will observe a rapid increase in the visible cross section of the Higgs production during energy scanning.

Figure 2 shows ten differential luminosities with the $J_z = 0$ angular momentum state of initial photon collisions in a photon-photon collider for energy scanning at $m_h = 100$ GeV. In this study, we have scanned the Higgs boson resonance from the left side to the right side in Fig. 2. The circles exhibit the luminosity points in contact with the Higgs boson and the rise of the luminosity at $m_h = 100$ GeV is rapid at the threshold of energy scanning. Here we introduce the required parameters for the luminosity calculation. A laser photon of energy ω_L is scattered by an electron beam of energy E_e in the conversion region of the photon-photon collider. The kinematics of Compton scattering is characterized by the dimensionless parameter [4]

$$x \equiv \frac{4E_e\omega_L}{m_e^2} \approx 15.3 \left[\frac{E_e}{\mathrm{TeV}}\right]\left[\frac{\omega_L}{\mathrm{eV}}\right], \tag{1}$$

where m_e is electron mass. The maximum energy of the scattered photon ω_{max} is $E_e x/(x+1)$ given by x. The parameter x is fixed to be 4.8, and we get $\omega_{max} = 100$ GeV when $E_e = 121$ GeV and $\omega_L = 2.6$ eV. The combination of the polarizations of the electron P_e and the laser P_L should be $P_L P_e = -1$ so that the generated photon spectrum peaks at its maximum energy. The differential luminosity distribution depends on the variable $\rho = b/(\gamma a)$, where a is the rms radius of the electron beam at the interaction point (IP), b is the distance between the conversion point (CP) and the IP, and $\gamma = E_e/m_e$. Polarized luminosities with the $J_z = 0$ and the $J_z = \pm 2$ in a photon-photon collider are described by

$$\frac{dL_{\gamma\gamma}^{J_z=0}}{dz} = 2zk^2 L_{ee} \int_{z^2/y_m}^{y_m} f(x,y) f(x,z^2/y) \times I_0\left(\rho^2\sqrt{\left(\frac{y_m}{y}-1\right)\left(\frac{y_m y}{z^2}-1\right)}\right) \times \exp\left[-\left(\frac{y_m}{y}+\frac{y_m y}{z^2}-2\right)\frac{\rho^2}{2}\right] \times \frac{1+P_1(y)P_2(z^2/y)}{2}\frac{dy}{y}, \tag{2}$$

$$\frac{dL_{\gamma\gamma}^{J_z=\pm 2}}{dz} = 2zk^2 L_{ee} \int_{z^2/y_m}^{y_m} f(x,y) f(x,z^2/y) \times I_0\left(\rho^2\sqrt{\left(\frac{y_m}{y}-1\right)\left(\frac{y_m y}{z^2}-1\right)}\right) \times \exp\left[-\left(\frac{y_m}{y}+\frac{y_m y}{z^2}-2\right)\frac{\rho^2}{2}\right] \times \frac{1-P_1(y)P_2(z^2/y)}{2}\frac{dy}{y}, \tag{3}$$

where $z = \sqrt{s}_{\gamma\gamma}/2E_e$, the center-of-mass energy of $\gamma\gamma$ collisions $\sqrt{s}_{\gamma\gamma}$, k the conversion coefficient, $y = \omega/E_e$, the energy of the scattered photon ω, $y_m = x/(x+1)$, L_{ee} the total luminosity of the e^+e^- collider, $f(y)$ the function of energy spectrum of the Compton scattering photons, $I_0(x)$ the modified Bessel function, $P_1(y), P_2(y)$ the polarization of left- and right-moving photons, respectively [4]. Here we assumed $\rho = 1$ and $k = 0.6$. It should be noted that the shape of the high energy edge and w_{max} are influenced by nonlinear effects due to very strong focus of the laser field at the CP. Prior to the actual energy scan, we need to have a fairly good estimate for nonlinear effects including the polarization.

Once the Higgs boson is observed at future e^+e^- colliders, we must determine its precise mass and width in order to reveal the fundamental properties of the Higgs boson. At a photon-photon collider, the feasibility of the measurement of the two-photon decay width of

a Higgs boson has been studied in the mass range $M_W < m_h < 2M_W$ [7–11]. For $m_h < 2M_W$, the SM Higgs boson mainly decays into a $b\bar{b}$ pair and the daughter b-flavored hadrons will be easily identified due to their long lifetime; therefore, the $b\bar{b}$ events are the best signals. The cross section of the Higgs boson production can be described by the Breit-Wigner approximation:

$$\sigma_{\gamma\gamma\to h\to b\bar{b}}(\sqrt{s}) = 8\pi\frac{\Gamma(h\to\gamma\gamma)\Gamma(h\to b\bar{b})}{(s-m_h^2)^2+m_h^2\Gamma_h^2}(1+\lambda_1\lambda_2), \tag{4}$$

where $\Gamma(h\to\gamma\gamma)$ and $\Gamma(h\to b\bar{b})$ are the decay widths of the Higgs boson into two photons and a $b\bar{b}$ pair, λ_1 and λ_2 the initial photon helicities, respectively. The effective cross section of the signal events within $m_h-\delta<\sqrt{s}<m_h+\delta$ is

$$\sigma^{\rm eff}_{\gamma\gamma\to h\to b\bar{b}} = \int_{m_h-\delta}^{m_h+\delta} 16\pi\frac{\Gamma(h\to\gamma\gamma)\Gamma(h\to b\bar{b})}{(\hat{s}-m_h^2)^2+m_h^2\Gamma_h^2}\frac{1}{L_{\gamma\gamma}}\frac{dL_{\gamma\gamma}^{J_z=0}}{d\sqrt{\hat{s}}}d\sqrt{\hat{s}}, \tag{5}$$

where δ expresses the effect of the detector resolution and we assumed $\delta = 5$ GeV. Here we supposed that the total luminosity is $L_{\gamma\gamma} = L_{\gamma\gamma}^{J_z=0} + L_{\gamma\gamma}^{J_z=\pm 2}$.

The main background processes may be the continuum $\gamma\gamma\to b\bar{b}$, $c\bar{c}$ as well as the radiative processes $\gamma\gamma\to b\bar{b}g$, $c\bar{c}g$. The continuum backgrounds dominantly produced by initial photon collisions in the $J_z=\pm 2$ can be suppressed by controlling the polarization of the colliding photon beams. Several authors reported that the effect of QCD corrections to $\gamma\gamma\to q\bar{q}$ is large since the helicity suppression which affects the background $q\bar{q}$ events does not work due to a gluon emission [9,10]. In this study we take account of the QCD corrections of the soft gluon emission, hard gluon emission, and virtual correction [10]. The effective cross section of the background process $\gamma\gamma\to b\bar{b}(g)$ or $c\bar{c}(g)$ within $m_h-\delta<\sqrt{s}<m_h+\delta$ is

$$\sigma^{\rm eff}_{\rm bg} = \int_{m_h-\delta}^{m_h+\delta}\sigma_{\rm bg}(\sqrt{\hat{s}})\frac{1}{L_{\gamma\gamma}}\frac{dL_{\gamma\gamma}}{d\sqrt{\hat{s}}}d\sqrt{\hat{s}}, \tag{6}$$

where $\sigma_{\rm bg}(\sqrt{s})$ is the cross section of the background events.

Since the cross section of $\gamma\gamma\to c\bar{c}$ is larger than that of $\gamma\gamma\to b\bar{b}$ due to the large electric charge of the quark, we apply the b tagging in order to eliminate the charm and the light

quark backgrounds. By the topological vertexing method [12] the $b\bar{b}(g)$ and $c\bar{c}(g)$ tagging efficiencies are assumed as 49% and 0.005%, respectively. We impose the following cuts to remove backgrounds: (1) the double $b\bar{b}$ tagging in the event; (2) $|\cos\theta_{b,\bar{b}}| < 0.95$, where $\theta_{b,\bar{b}}$ is the scattering angle of the $b(\bar{b})$ quark; (3) $|M_{b\bar{b}} - m_h| < 5$ GeV.

Figure 3 shows an example of energy scan to determine m_h. Each energy point corresponds to 5 fb^{-1} and the total luminosity of photon-photon collisions is 50 fb^{-1} in the same distributions as with Fig. 2. The total width of the SM Higgs boson $\Gamma_{h\ \mathrm{SM}}$ for m_h=100 GeV is 2.16 MeV, which is computed by the HDECAY program [13]. The partial widths $\Gamma(h \to \gamma\gamma)$ and $\Gamma(h \to b\bar{b})$ at the predicted SM value at $m_h = 100$ GeV are fixed for energy scanning at $m_h = 99.8$, 100, 100.2 GeV. The statistical errors in Fig. 3 indicate $\sqrt{S+B}$, where S and B are the number of signal and background. From figure 3, we can understand that the method of energy scanning for m_h is more effective than that of the measurement of a single point at the luminosity peak using the same total luminosity, because the statistical errors at the threshold of energy scanning are smaller than that at the luminosity peak and we can distinguish the mass difference of 200 MeV. With the energy scanning of 10 points, the attainable error in m_h is about 110 MeV at the 1σ level.

The measurement for the determination of Γ_h by the method of energy scan is shown in Fig. 4. Each energy point corresponds to 5 fb^{-1} and the total luminosity is 50 fb^{-1}. The partial widths $\Gamma(h \to \gamma\gamma)$ and $\Gamma(h \to b\bar{b})$ at the predicted SM value at $m_h = 100$ GeV are fixed for energy scanning $\Gamma_h/\Gamma_{h\ \mathrm{SM}} = 0.9$, 1.0, 1.1. The large difference between the total widths at the luminosity peak can be seen easily in Fig. 4. The statistical error in Γ_h is about 6% at the 1σ level. If there are additional invisible decay modes of Higgs boson, only the total decay width increases keeping the partial widths of two photons and a $b\bar{b}$ pair unchanged. It is in this study of $\Gamma_h/\Gamma_{h_{\mathrm{SM}}} > 1$. Of course, this deviation from the SM should have also been observed in the parent e^+e^- collider. However, it is an independent observation in gamma-gamma energy scan which conforms the e^+e^- result.

Here we consider two cases of photon-photon colliders. First case is to be adopt at $x = 4.8$ and the energies of laser photon and electron beam are tuned in the case during

energy scan. Second case is to be choosen the fixed laser energy and the energy of electron beam is tuned only during energy scan. First and second cases are named tunable case and fixed case, respectively.

Table I lists the statistical errors of the SM Higgs boson mass at the 1σ level, using an integrated luminosity of 50/10 fb^{-1}. In this table, the mass errors of the tunable case are almost smaller than those of the fixed case. Since the background processes $\gamma\gamma \to q\bar{q}(g)$ are increasing at the lower Higgs mass and the branching ratio $B(h \to b\bar{b})$ is decreasing at the higher Higgs mass, the errors of the Higgs boson mass near 100 GeV are smallest. The statistical errors $\sqrt{S+B}/S$ of the total width $\Gamma_h/\Gamma_{h\ \mathrm{SM}}$ of the SM Higgs boson with a 50 fb^{-1} luminosity are listed in Table II. The statistical errors of the total width for intermediate-mass Higgs bosons are almost within 8% in Table II.

As for other origins of the errors, we need to know the systematic uncertainties on the luminosity distribution. The possibilities of the luminosity measurements at photon-photon colliders have been studied using the process $\gamma\gamma \to l^+l^-$ or $\gamma\gamma \to W^+W^-$ [14,15]. For energy scanning the measurement of the luminosity distribution at the high energy-edge is crucial and we need to study it further.

In summary, we have shown that it is possible to determine the Higgs boson mass to a high precision by the method of energy scanning at photon-photon colliders, using the high energy edge of the photon spectrum. The measurement of the total width of the Higgs boson at photon-photon colliders could distinguish a standard model Higgs boson from the extended models such as the minimal supersymmetric model when the partial widths $\Gamma(h \to \gamma\gamma)$ and $\Gamma(h \to b\bar{b})$ are fixed at the predicted standard model value.

I express sincere thanks to T. Takahashi, T. Tauchi, and I. Watanabe for useful discussions. This work was supported in part by the U.S. Department of Energy under Contract No. DE-AC03-76SF00098.

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TABLES

TABLE I. The statistical errors of the SM Higgs boson mass by devoting 50/10 fb^{-1} to each point.

			$\Delta m_{h_{\rm SM}}$ (MeV)			
$m_{h_{\rm SM}}$ (GeV)	80	90	100	110	120	140
Tunable case	$+140$	$+120$	$+100$	$+140$	$+170$	$+210$
	-150	-140	-110	-100	-100	-220
Fixed case	$+140$	$+120$	$+100$	$+170$	$+200$	$+220$
	-170	-130	-120	-110	-130	-270

TABLE II. The statistical errors of the total width of the SM Higgs boson by devoting 50/10 fb^{-1} to each point.

			$\Delta\Gamma_h/\Gamma_{h_{\rm SM}}$ (%)			
$m_{h_{\rm SM}}$ (GeV)	80	90	100	110	120	140
Tunable case	$+7.8$	$+6.7$	$+6.0$	$+5.7$	$+5.7$	$+7.5$
	-7.0	-6.1	-5.6	-5.3	-5.6	-6.8
Fixed case	$+8.5$	$+7.3$	$+6.6$	$+6.2$	$+6.3$	$+8.1$
	-7.6	-6.6	-6.0	-5.7	-5.8	-7.3

FIGURES

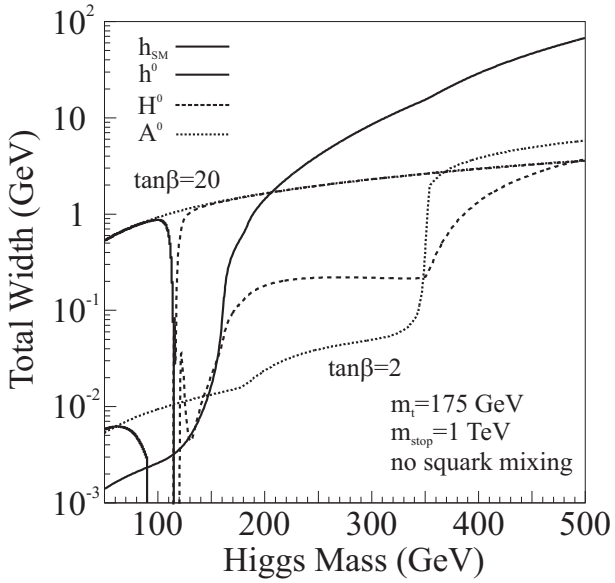

FIG. 1. The total widths of the SM and MSSM Higgs bosons. The top quark mass is assumed to be 175 GeV. In the case of the MSSM, the results for $\tan\beta = 2$ and 20 are shown, taking $m_{\tilde{t}} = 1$ TeV, including two-loop radiative corrections, and neglecting squark mixing. SUSY decay channels are assumed to be absent. Computed by HDECAY [13].

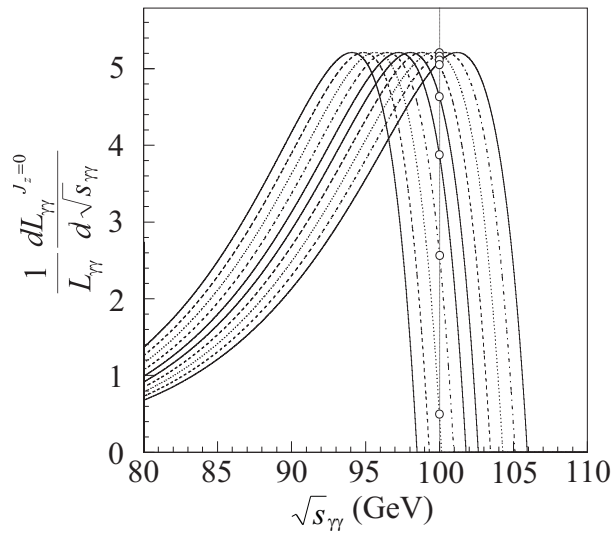

FIG. 2. Ten differential luminosities with $J_z = 0$ as a function of the center-of-mass energy in a photon-photon collider for energy scanning at $m_h = 100$ GeV.

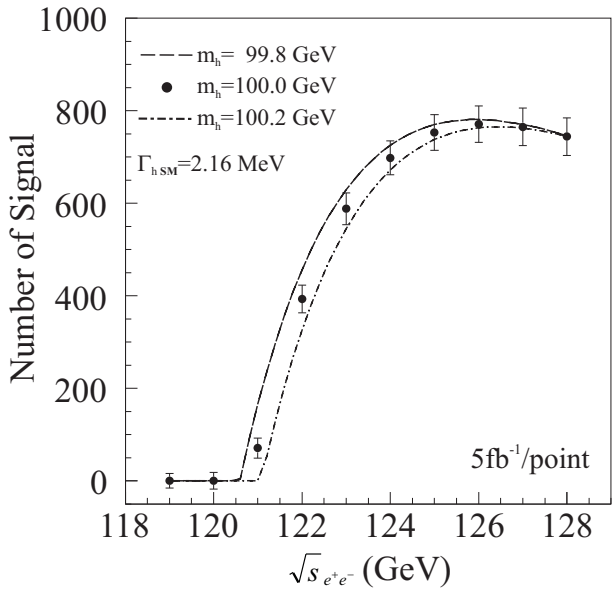

FIG. 3. An example of energy scan to determine m_h where each point corresponds to 5 fb^{-1}.

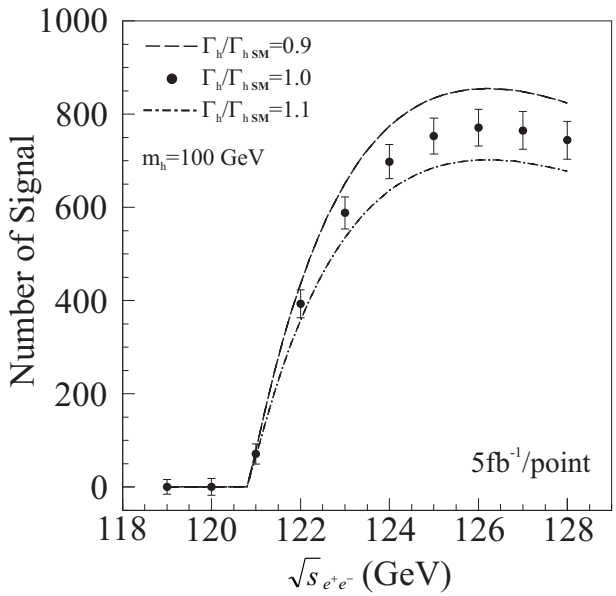

FIG. 4. An example of energy scan to determine Γ_h where each point corresponds to 5 fb^{-1}.