



Searching for New Physics in B decays

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New era of B physics

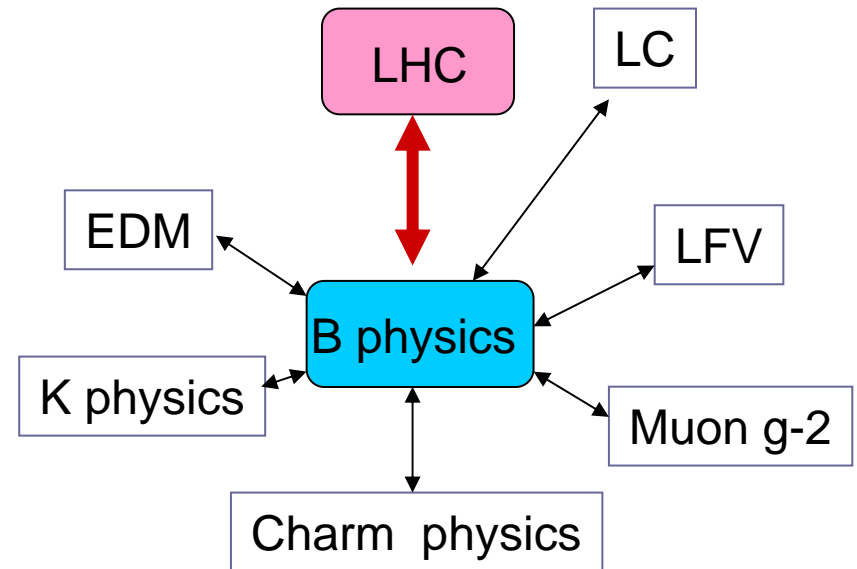
- Two B factory experiments Belle at KEKB and BABAR at PEP-II are very successful.
(~340/fb at KEKB and ~250/fb at PEP-II)
- The asymmetric B factories provides measurements of time-dependent CP violations in B decays.
- In future, more B physics will come at hadron machines (Tevatron, LHCb) and upgrade of the current B factories as well as Super B Factory (5-10/ab/year).

Goals of future B physics

- Main purpose of B physics from now on is to search for new physics effects in flavor-mixing and CP violation.
- There are several ways to look for new physics in CP violation and rare B decay processes.
- In order to identify a new physics model, we need to know pattern of deviations from the SM predictions in various observables.

New Physics in LHC era

- Some signals of new physics may be obtained at early stage of LHC.
(SUSY, Large extra dim. etc)
- Important to consider impacts of B physics to LHC physics, and vice versa.
- In general, correlations among various areas are important to figure out what is new physics.



Content of this talk

- Various methods to look for new physics effects in B decays.
- Comparative study of B physics signals in three SUSY models. “SUSY loop effects”
- SUSY with minimal flavor violation (MFV) at a large $\tan \beta$. “Higgs exchange”
- B physics signals of large extra dimension models

Super KEKB Lol

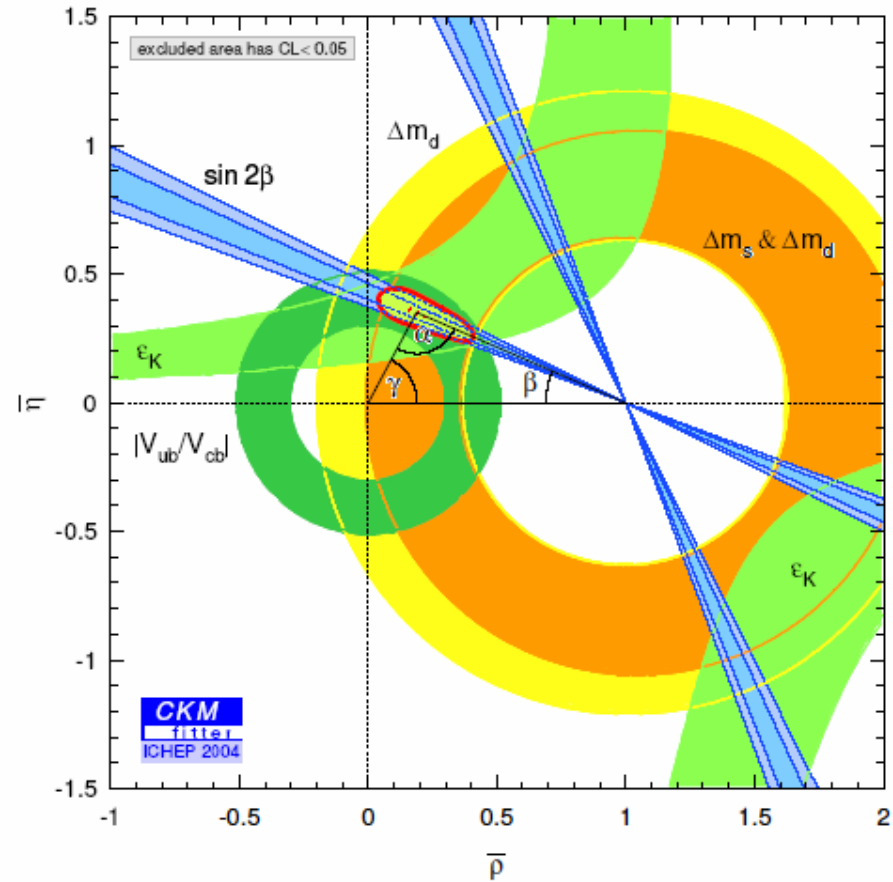
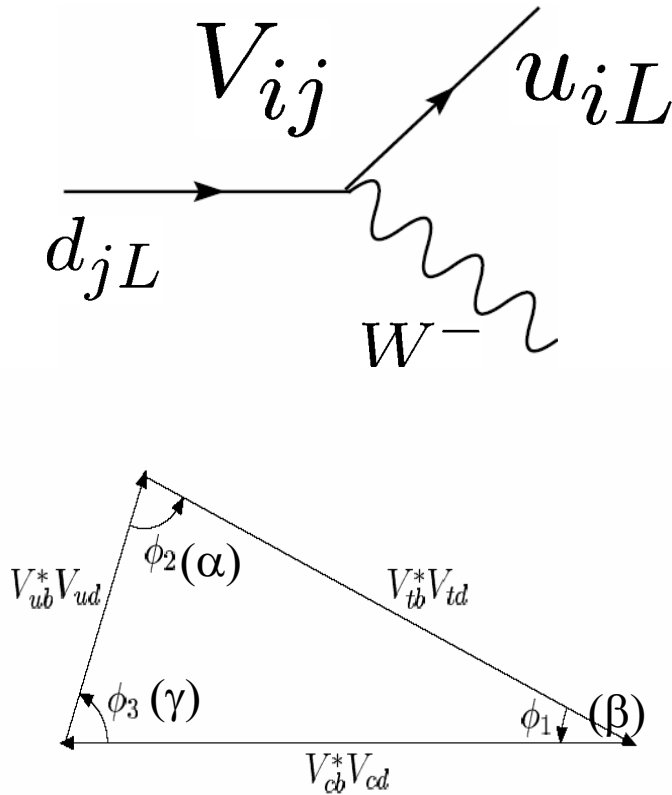
SLAC 10³⁶ study group

New Physics Searches in B decays

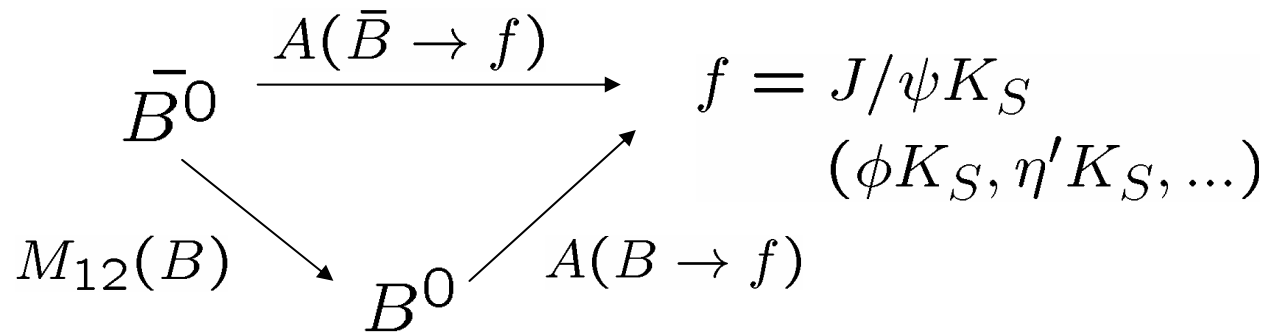
- There are many ways to look for new physics effects in B decays, both at Bd(Bu) and Bs experiments.
 1. Consistency test of the unitarity triangle.
 2. Comparison of various CP asymmetries.
 3. Rare B decays.
 4. Tau and charm physics at a Super B factory.

Unitarity triangle

- In the SM, all flavor and CP phenomena can be explained by the Cabibbo- Kobayashi-Maskawa matrix.



Time-dependent CP asymmetry in B decays



Time –dependent asymmetry can arise from the interference of two paths in the $B \rightarrow f$ decay amplitude.

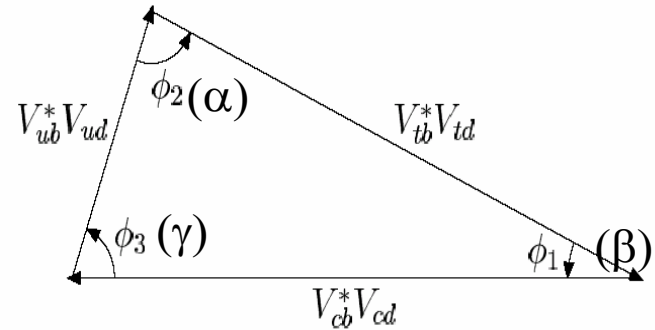
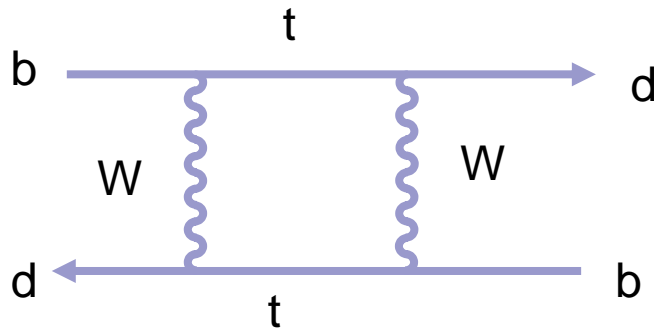
$$\frac{\Gamma(\bar{B}_{phys}^0 \rightarrow f) - \Gamma(B_{phys}^0 \rightarrow f)}{\Gamma(\bar{B}_{phys}^0 \rightarrow f) + \Gamma(B_{phys}^0 \rightarrow f)} = A_f \cos \Delta mt + S_f \sin \Delta mt$$

A_f : Direct CP asymmetry

S_f : Mixing –induced (Time-dependent) CP asymmetry

In the Standard Model, the B- B bar mixing amplitude have the phase $2\phi_1$

$$M_{12} \propto (V_{td}^* V_{tb})^2 \propto e^{2i\phi_1}$$



In general, the decay amplitude depends on several weak phases

$$A(B \rightarrow f) = \sum_k a_k e^{i\delta_k} e^{i\phi_k^W}$$

$$A(\bar{B} \rightarrow f) = \eta_f \sum_k a_k e^{i\delta_k} e^{-i\phi_k^W} \quad (\eta_f : \text{CP eigenvalue of } f)$$

If f is an CP eigen state, and the decay amplitude is dominated by one weak phase amplitude, then

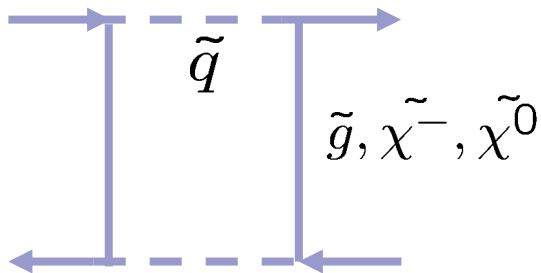
$$A_f = 0, S_f = -\eta_f \sin 2(\phi_1 + \phi^W)$$

New physics effects in B-B mixing

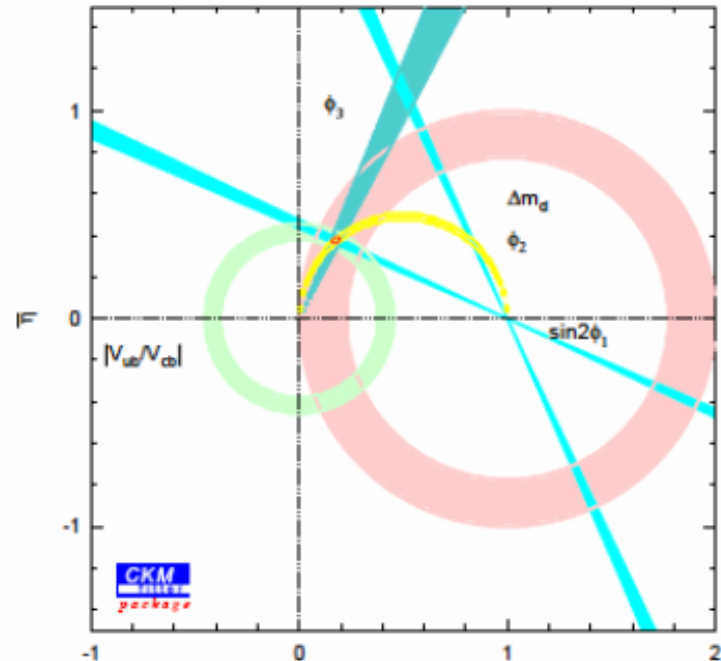
1. Bd-unitarity triangle

→ New contributions to the Bd mixing amplitude.

Ex. SUSY loop diagram



Super KEKB Lol ,hep-ex/0406071



Unitarity triangle at 50/ab

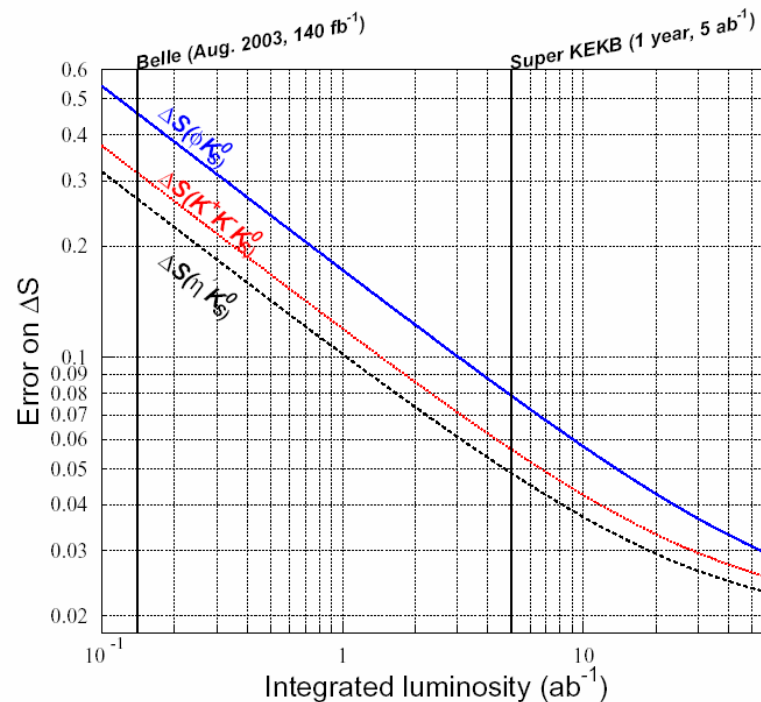
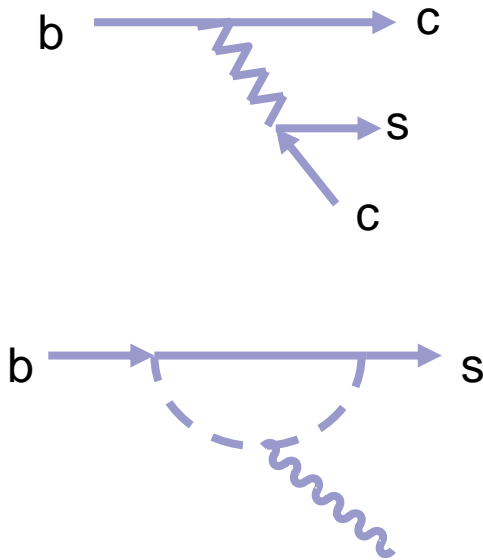
2. Bs mixing and CP asymmetry in $B_s \rightarrow J/\psi \phi$

→ The magnitude and the phase of the Bs mixing amplitude.

Comparison of various CP asymmetries

CP asymmetries from $B \rightarrow J/\psi K_s$, $B \rightarrow \phi K_s$, $B \rightarrow \eta' K_s$.

→ A new CP phase in the b-s-g amplitude. These should be the same in the SM.



In order to confirm the anomaly of b-s transition, we need a large luminosity (>a few /ab)

Rare B decays

Direct CP violation in $b \rightarrow s \gamma$. (New phase in $b \rightarrow s \gamma$)

$$A_{CP} = \frac{B(b \rightarrow s \gamma) - B(\bar{b} \rightarrow s \gamma)}{B(b \rightarrow s \gamma) + B(\bar{b} \rightarrow s \gamma)} \quad |A_{CP}| < 1\% \text{ in SM}$$

Mixing induced CP violation in $B \rightarrow M s \gamma$. ($b \rightarrow s \gamma_R$)

$$A_{CP}^{mix}(B \rightarrow K^* \gamma) = \frac{2 \text{Im}(e^{-i\phi_M} C_7 C_7')}{|C_7|^2 + |C_7'|^2}$$

$$H = \frac{4G_F}{\sqrt{2}} \{C_7'(\bar{s}_R \sigma^{\mu\nu} b_L)F_{\mu\nu} + C_7(\bar{s}_L \sigma^{\mu\nu} b_R)F_{\mu\nu} + h.c\} \quad A_{CP} \sim O(\text{ms/mb}) \text{ in SM}$$

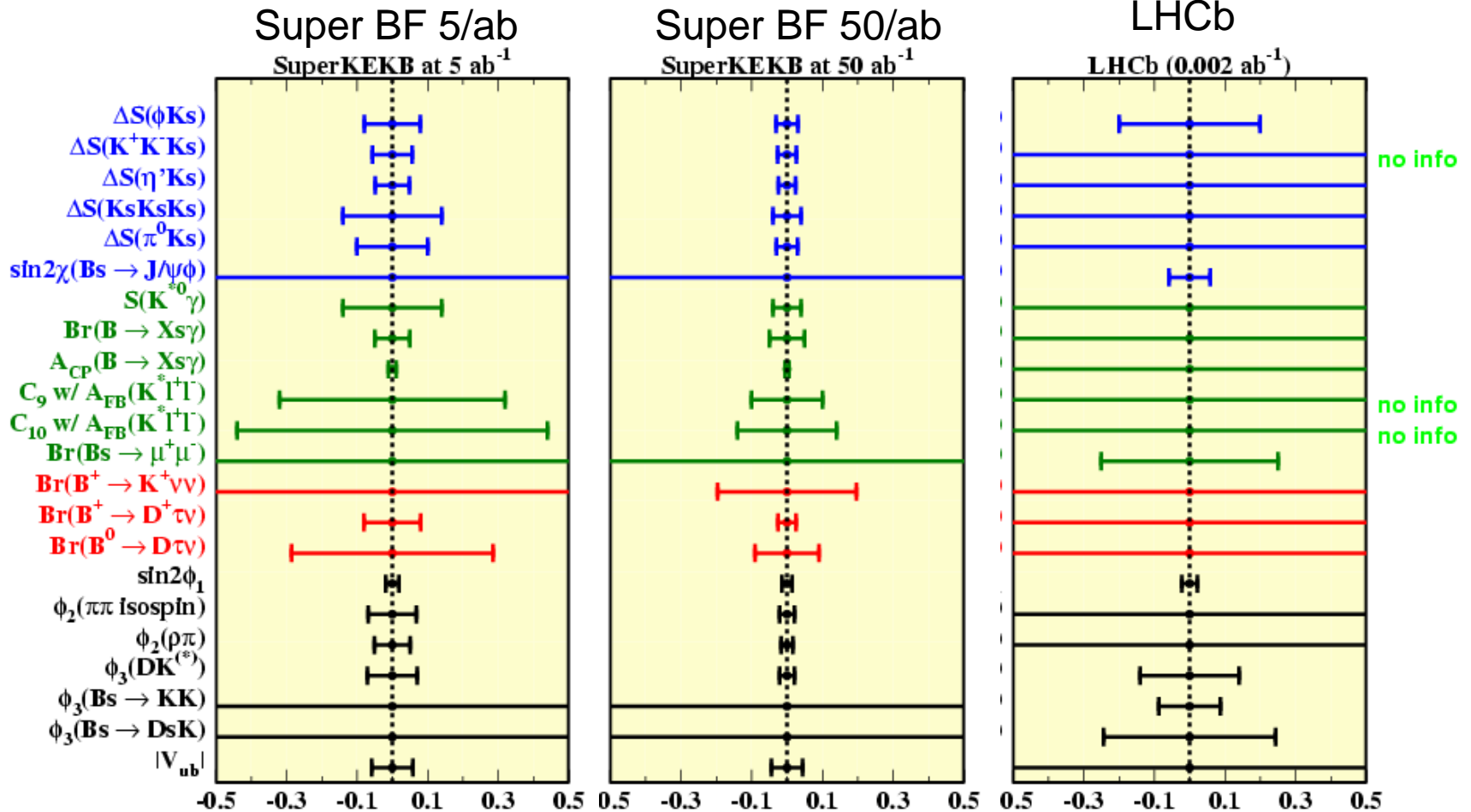
Branching ratio and lepton FB asymmetry in $b \rightarrow s \ell \ell$.

$b \rightarrow s \nu \nu$, $B \rightarrow \tau \tau$, $B \rightarrow \ell \ell$,

$B \rightarrow D \tau \nu$ (Charged Higgs exchange)

Summary of physics reach

Super KEKB Lol



Complementarity between Super B Factory and hadron B programs

SUSY and Flavor Physics

- SUSY modes introduce SUSY partners.
- Squark mass matrixes are new sources of flavor mixing and CP violation.
- Squark masses depend on SUSY breaking terms as well as the Yukawa coupling constants.

quark (q)
lepton (l)
gluon (g)
W,Z, γ ,H

squark (\tilde{q})
slepton (\tilde{l})
gluino (\tilde{g})
neutralino, chargino ($\tilde{\chi}$)

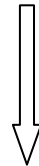
Quark mass $(m_q)_{ij} = Y_{ij}v$

Squark mass $(m_{\tilde{q}}^2)_{ij} = (Y^\dagger Y)_{ij}v^2 + m_{ij}^2$

SUSY breaking

- Squark mass matrixes carry information on the SUSY breaking mechanism and interactions at the GUT scale.

Origin of SUSY breaking
(mSUGRA, AMSB, GMSB,
Flavor symmetry, etc.)



← Renormalization
(SUSY GUT, neutrino Yukawa couplings etc.)

SUSY breaking terms at the M_w scale
(squark, slepton, chargino, neutralino, gluino masses)

Diagonal : LHC/LC
Off-diagonal: Future Flavor exp.



Top quark: Tevatron
KM phase: B factories

$$(m_{\tilde{q}}^2)_{ij} = \begin{pmatrix} m_{11}^2 & m_{12}^2 & m_{13}^2 \\ m_{21}^2 & m_{22}^2 & m_{23}^2 \\ m_{31}^2 & m_{32}^2 & m_{33}^2 \end{pmatrix}$$

Distinguishing different SUSY models

T.Goto, Y.Okada, Y.Shimizu, T.Shindou, and M.Tanaka

- In order to illustrate the potential of B physics in exploring flavor structure of SUSY breaking, we calculate various observables in four cases of SUSY models.

Models

1. Minimal supergravity model
2. SU(5) SUSY GUT with right-handed neutrino
 - 2-1. degenerate RHN case
 - 2-2. non-degenerate RHN case
3. MSSM with U(2) flavor symmetry

Observables

- Bd-Bd mixing, Bs-Bs mixing.
- CP violation in K-K mixing (ϵ).
- Time-dependent CP violation in $B \rightarrow J/\psi K_s$, $B \rightarrow \phi K_s$, $B \rightarrow K^* \gamma$.
- Direct CP violation in $b \rightarrow s \gamma$.

Three SUSY Models

Origin of the squark mixing

$$(m_{\tilde{q}}^2)_{ij} = (Y^\dagger Y)_{ij} v^2 + m_{ij}^2$$

1. Minimal supergravity model. Only the CKM matrix

Minimal Flavor Violation

2. SU(5) SUSY GUT with right-handed neutrino.

Neutrino Flavor Mixing

The CKM matrix and the neutrino Yukawa coupling constants

2-1. degenerate RHN case ($\mu \rightarrow e \gamma$ large)

2-2. non-degenerate case ($\mu \rightarrow e \gamma$ suppressed)

$$m_{ij}^2 \sim c (y_\nu^\dagger y_\nu)_{ij}$$

3. MSSM with U(2) flavor symmetry.

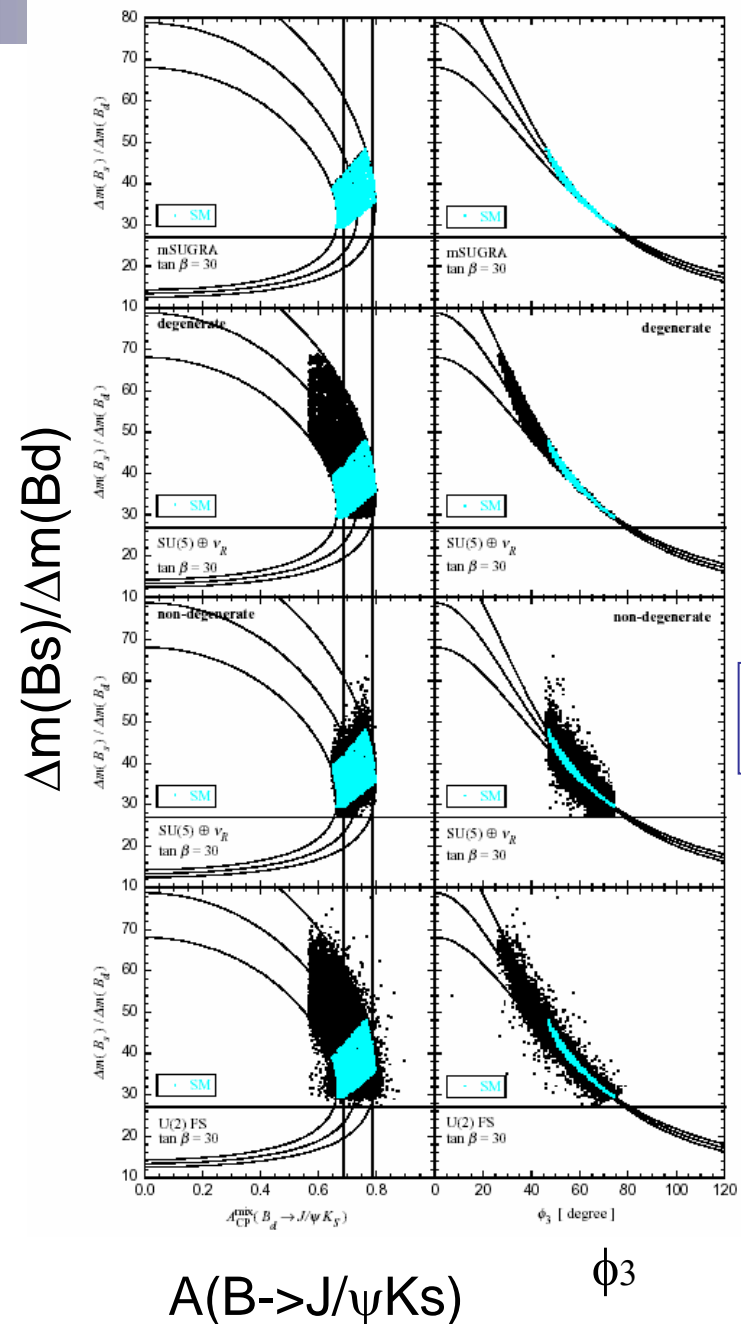
Both Yukawa coupling constants and SUSY breaking terms have the (12)-3 structure.

$$(m_{\tilde{q}}^2)_{ij} \simeq \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 + O(\epsilon^2) & O(\epsilon) \\ 0 & O(\epsilon) & O(1) \end{pmatrix}$$

Approximate Flavor Symmetry

Unitarity triangle

- Small deviation in mSUGRA.
- B_d unitarity triangle is closed, but ε_K has a large SUSY contribution in SU(5) GUT for the degenerate M_R case.
- B_s mixing receives SUSY effects for the non-degenerate case.
- Various SUSY contributions for the U(2) flavor symmetry model.



mSUGRA

SU(5) GUT
Degenerate

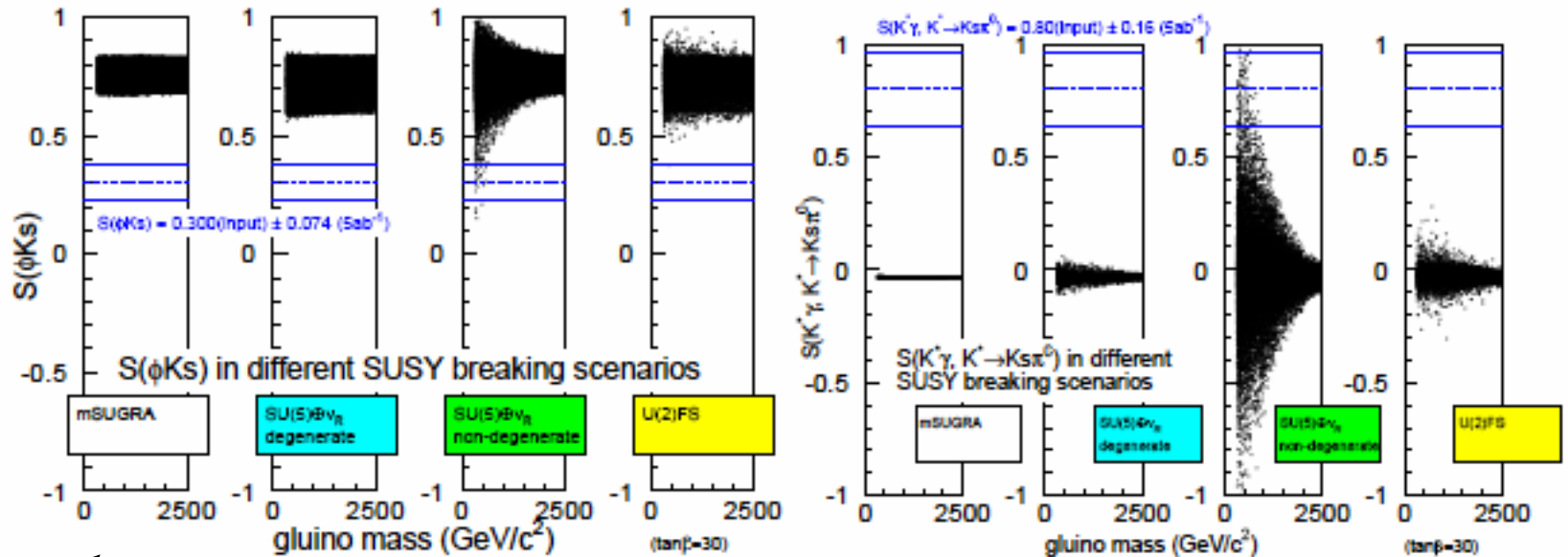
SU(5) GUT
Non-degenerate

U(2) FS

$A(B \rightarrow J/\psi K_S)$

ϕ_3

CP asymmetries in $B \rightarrow \phi K_s$ and $b \rightarrow s \gamma$



CP asymmetry
in $B \rightarrow \phi K_s$

Direct asymmetry
in $b \rightarrow s \gamma$

CP asymmetry
in $B \rightarrow K^* \gamma$

In the SU(5) +RHN with non-degenerate Majorana masses, if $S(\phi K_S) < 0.5$, then

Collider physics

Gluino mass $< 800 \text{ GeV}$ } LHC
 Stop mass $< 800 \text{ GeV}$ }
 Lighter chargino mass $< 300 \text{ GeV}$ } LC
 2nd neutralino mass $< 300 \text{ GeV}$ }

LFV

$$B(\mu \rightarrow e\gamma) > 10^{-14}$$

$$B(\tau \rightarrow \mu\gamma) > 10^{-9}$$

EDM

$$\text{neutron EDM} > 10^{-26}$$

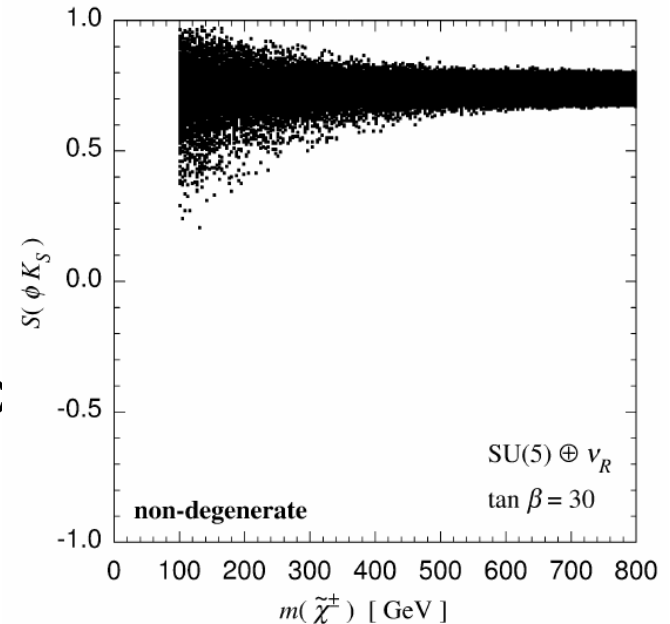
$$\text{Hg EDM} > 10^{-29}$$

MEG (PSI) , MECO (BNL, μ -e conv. exp)

Super B factory

One order of magnitude below the current bounds

$S(\phi K_S)$ vs chargino mass



Pattern of deviations from the SM prediction

	B_d unitarity			A_{CP}^{mix} $B \rightarrow \phi K_S$	A_{CP}^{mix} $B \rightarrow K^* \gamma$	A_{CP}^{dir} $B \rightarrow X_s \gamma$	A_{CP}^{mix} $B_s \rightarrow J/\psi \phi$
	closure	$+\epsilon_K$	$+\Delta m(B_s)$				
mSUGRA	closed	-	-	-	-	-	-
SU(5) SUSY GUT (degenerate RHN)	closed	✓	-	-	-	-	-
SU(5) SUSY GUT (non-deg. RHN)	closed	-	✓	✓	✓	-	✓
MSSM with U(2)	✓	✓	✓	✓	✓	✓	✓

mSUGRA: small deviation

SUSY SU(5) with degenerate RHN: signals in 1-2 mixing

SUSY SU(5) with non-degenerate RHN: signals in 2-3 mixing

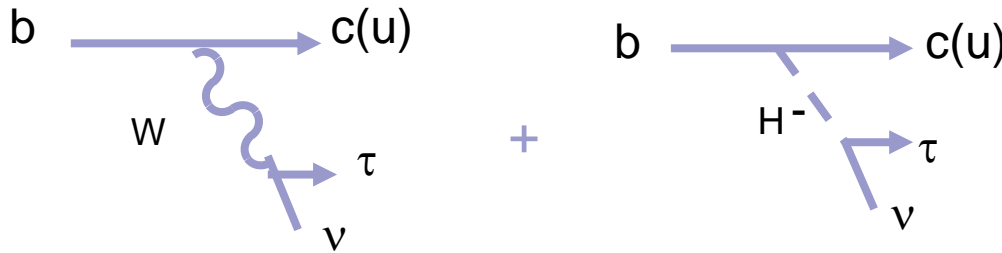
MSSM with U(2) FS: various new physics signals

SUSY with a minimal flavor violation (MFV)

- Even in the case where the squark flavor mixing is similar to the quark flavor mixing (MFV), a large deviation from the SM is possible for a large value of two vacuum expectation values ($\tan \beta$).
- Effects can be significant for the charged Higgs boson exchange in $B \rightarrow D \tau \nu$ and $B \rightarrow \tau \nu$.
- $B_s \rightarrow \mu \mu$ is enhanced by the loop-induced flavor changing neutral Higgs coupling.

Tauonic B decay, $B \rightarrow D \tau \nu$, $B \rightarrow \tau \nu$

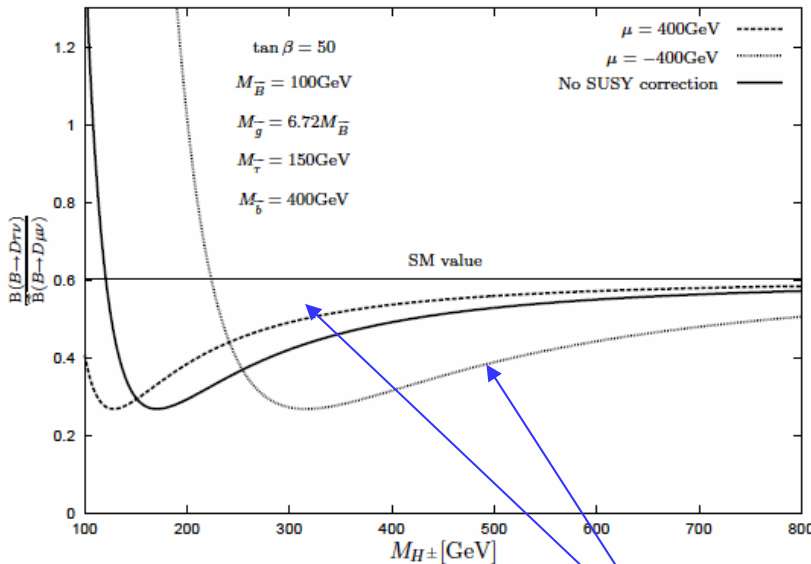
Charged Higgs boson exchange.



$B \rightarrow D \tau \nu$

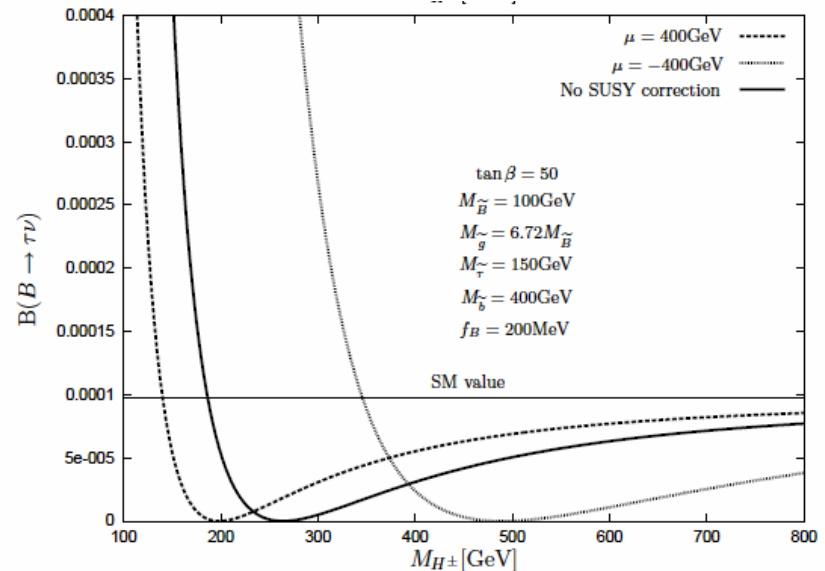
$\tan \beta = 50$

$B \rightarrow \tau \nu$



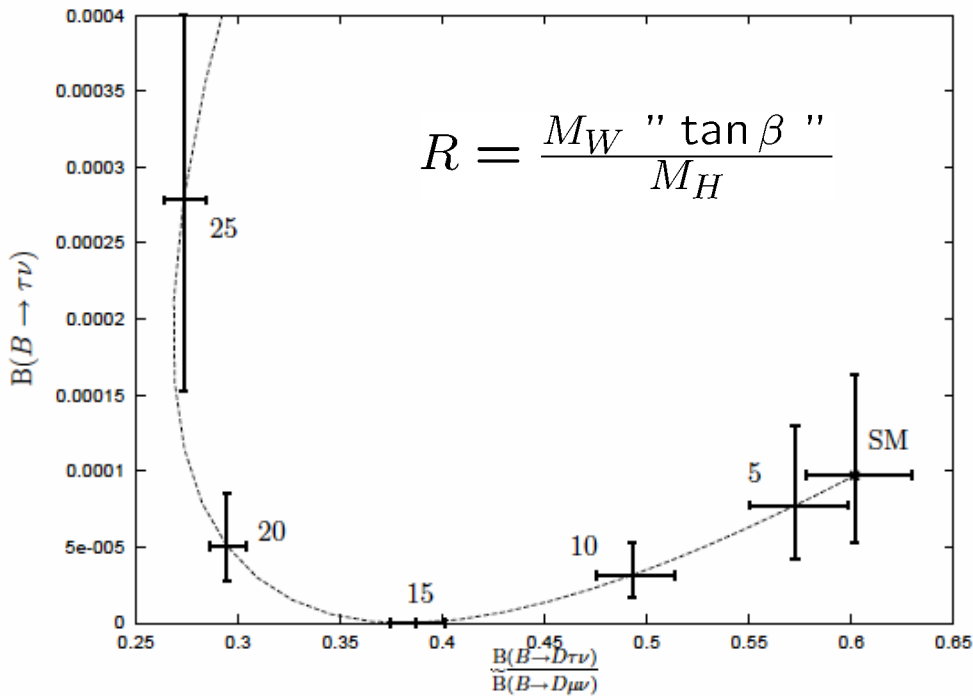
Charged Higgs mass

SUSY loop corrections
to the Higgs vertex



Charged Higgs mass

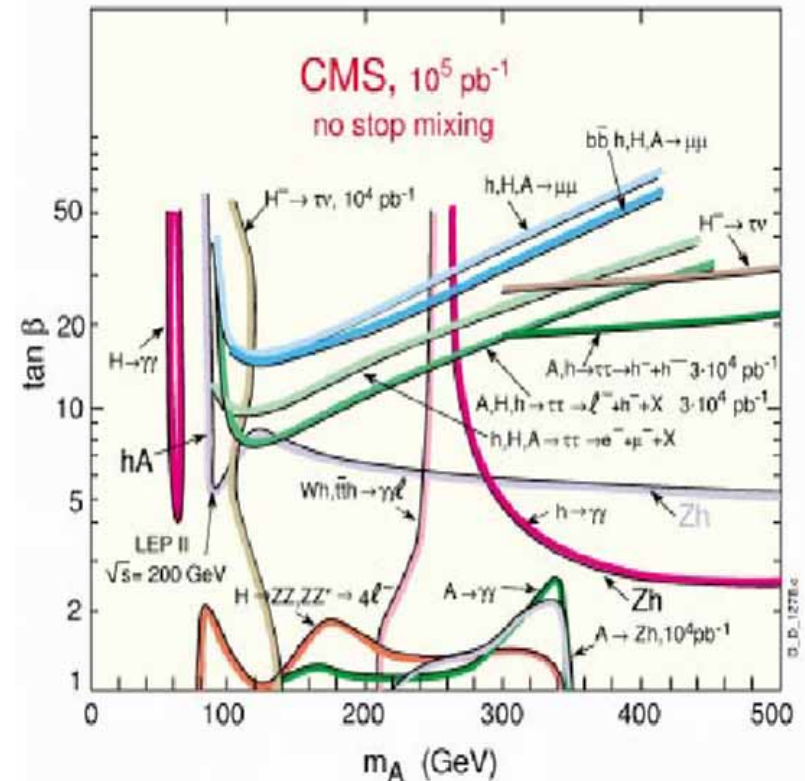
Correlation between $B \rightarrow D \tau \nu$, $B \rightarrow \tau \nu$



Super KEKB sensitivity from $B \rightarrow D \tau \nu$

$$M_H > M_W \tan \beta / 11 \text{ @ } 5/\text{fb (90\%CL limit)}$$

LHC heavy Higgs boson search

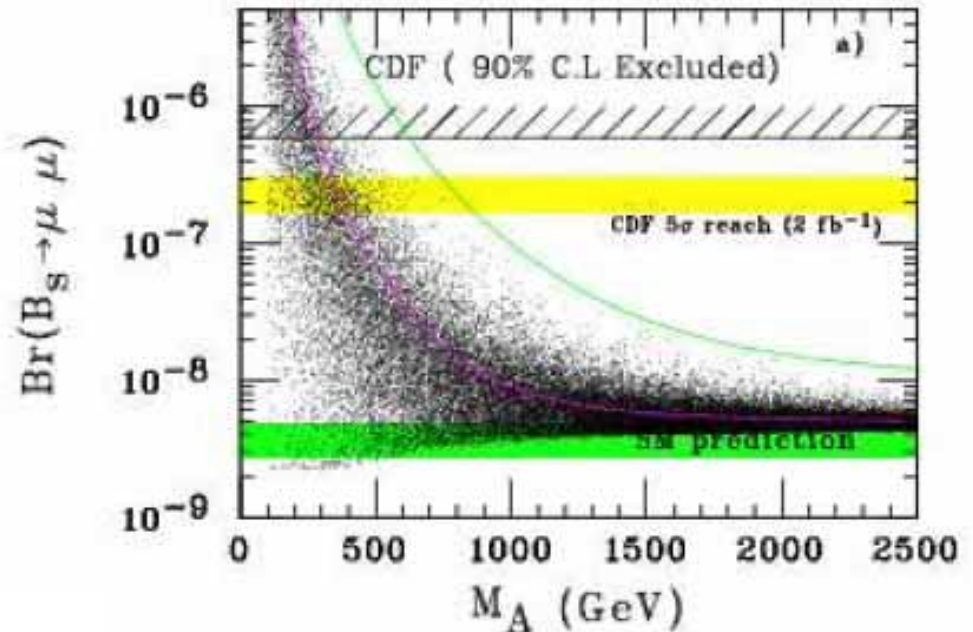
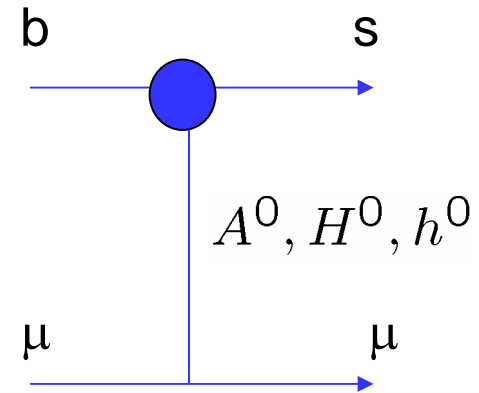


The covered parameter space
Is similar to LHC direct Higgs search

$B(B_s \rightarrow \mu\mu)$

Loop-induced neutral Higgs exchange effects

- SUSY loop corrections can enhance $B(B_s \rightarrow \mu\mu)$ by a few orders of magnitude from the SM prediction for large values of $\tan\beta$. This is within the reach of Tevatron exp.

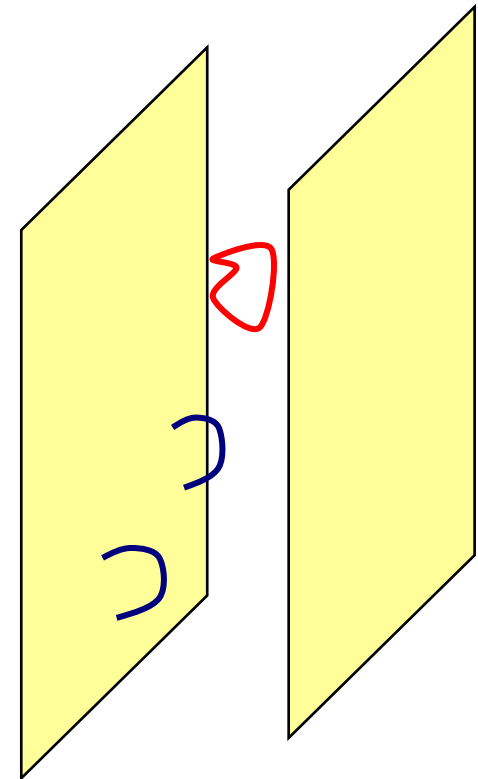


A.Dedes, B.T.Huffman

$$B(B_s \rightarrow \mu\mu) \sim 5 \times 10^{-7} \left(\frac{\tan\beta}{50}\right)^6 \left(\frac{300\text{GeV}}{M_A}\right)^4$$

Large extra dim and B physics

- Models with large extra dimensions were proposed as an alternative scenario for a solution to the hierarchy problem.
- Various types of models:
Flat extra dim vs. Curved extra dim
What particles can propagate in the bulk.
- Geometrical construction of the fermion mass hierarchy
=> non-universality of KK graviton/gauge boson couplings



KK graviton exchange

KK graviton exchange can induce tree-level FCNC coupling.

$$O_{grav} = \frac{X}{M^4} T_{\mu\nu} T^{\mu\nu}$$

Differential branching ratio of $b \rightarrow sll$ processes.

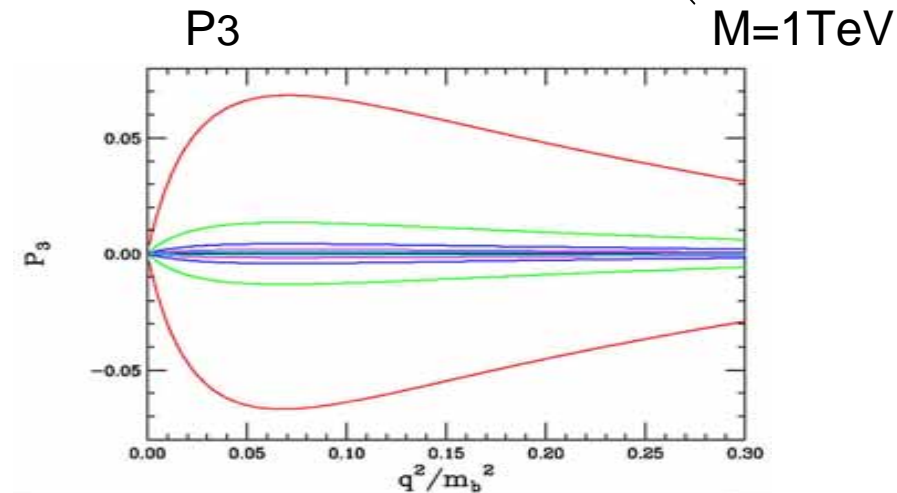
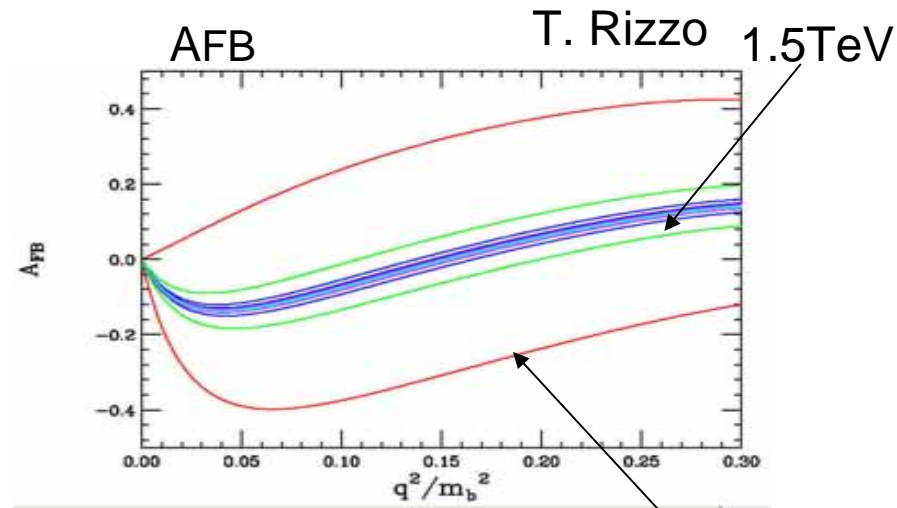
P3 : 3rd Legendre polynomial moment

=> pick up $(\cos\theta)^3$ terms due to spin2 graviton exchange.

(In both flat and curved extra dim)

T.Rizzo

b->sll differential Br



(Flat large extra dim case)

KK gluon, KK Z-boson exchange in warped extra dim.

In the warped extra dimension with bulk fermion/gauge boson propagation in order for the fermion mass hierarchy, we put

Light fermion \rightarrow localized toward Planck brane

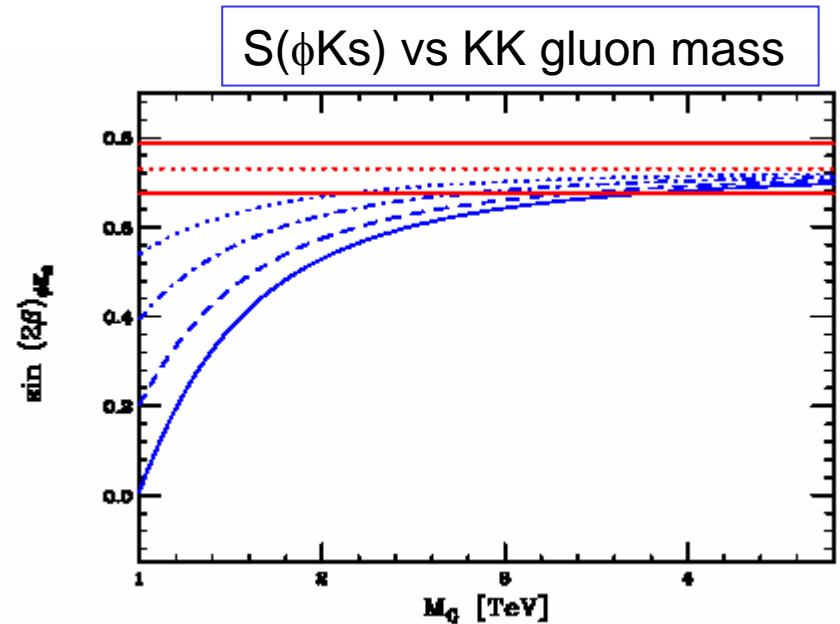
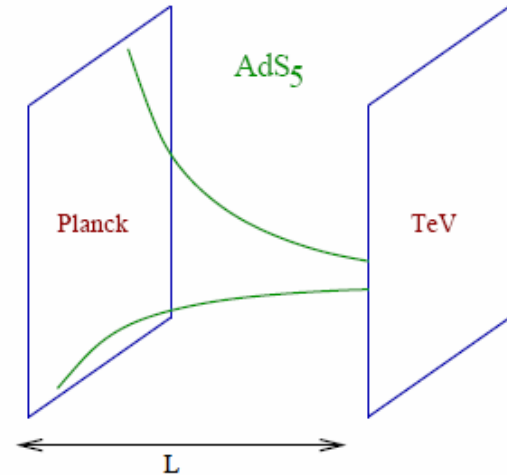
Top and left-handed bottom \rightarrow localized toward the TeV brane.

\Rightarrow Generate tree level FCNC in KK gluon and Z boson exchange.

Various FCNC four fermion interactions

($S(\phi K_s)$, $b \rightarrow sll$, B_s -mixing, etc.)

A. Agashe, et.al; G.Burdman



1st KK gluon mass

Summary of new physics signals

Model	B_d Unitarity	Time-dep. CPV	Rare B decay	Other signals
mSUGRA(moderate $\tan \beta$)	-	-	-	-
mSUGRA(large $\tan \beta$)	B_d mixing	-	$B \rightarrow (D)\tau\nu$ $b \rightarrow sl^+\ell^-$	$B_s \rightarrow \mu\mu$ B_s mixing
SUSY GUT with ν_R	-	$B \rightarrow \phi K_S$ $B \rightarrow K^*\gamma$	-	B_s mixing τ LFV, n EDM
Effective SUSY	B_d mixing	$B \rightarrow \phi K_S$	$A_{CP}^{b \rightarrow s\gamma}, b \rightarrow sl^+\ell^-$	B_s mixing
KK graviton exchange	-	-	$b \rightarrow sl^+\ell^-$	-
Split fermions in large extra dimensions	B_d mixing	-	$b \rightarrow sl^+\ell^-$	$K^0\bar{K}^0$ mixing $D^0\bar{D}^0$ mixing
Bulk fermions in warped extra dimensions	B_d mixing	$B \rightarrow \phi K_S$	$b \rightarrow sl^+\ell^-$	B_s mixing $D^0\bar{D}^0$ mixing
Universal extra dimensions	-	-	$b \rightarrow sl^+\ell^-$ $b \rightarrow s\gamma$	$K \rightarrow \pi\nu\bar{\nu}$

In “The Discovery Potential of a Super B Factory”,
The Proceedings of the 2003 SLAC Workshops

Summary

- Flavor physics tell us important aspects of new physics models.

SUSY -> interactions at high energy scale.

Large Extra Dim -> origin of fermion mass/flavor structure.

- There are a variety of ways to look for new physics effects in B decays.
- In order to distinguish different models, we need to know the pattern of deviations from SM predictions.
- Mutual impacts among B physics, K/D physics, LHC/LC, LFV, EDM, etc. are important.