

# Lectures on Brane World Scenarios V

Today's topics:

“Diluting the cosmological constant in infinite volume extra dimensions”

by G. Dvali, G. Gobadadze and M. Shifman

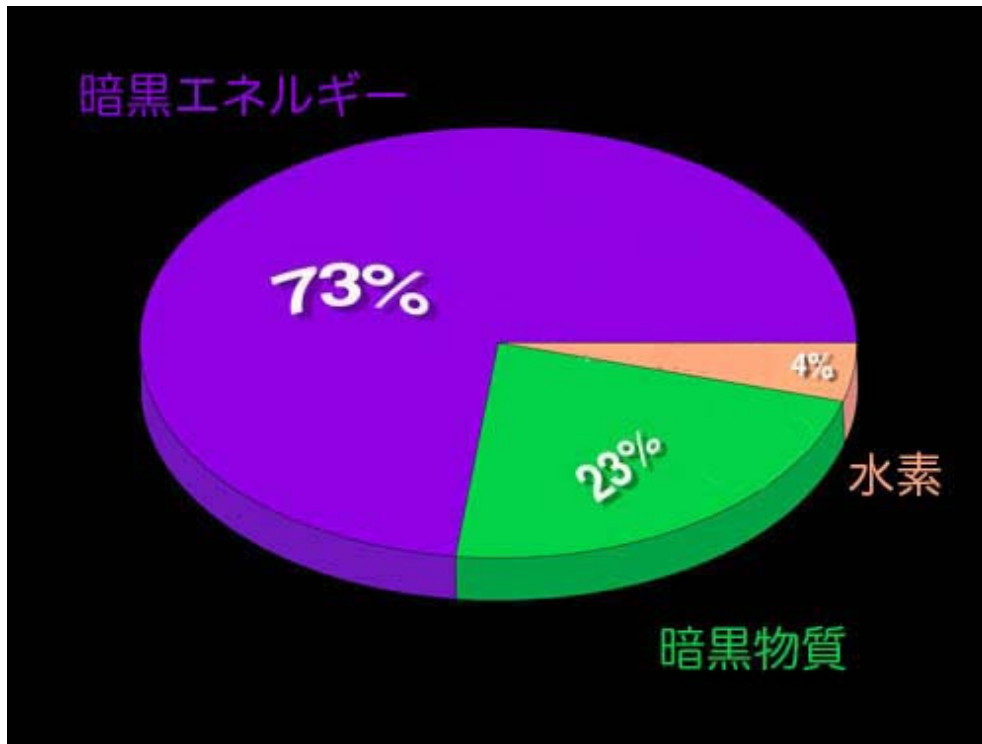
PRD 67, 044020 (2003) , and references therein & thereof

**Nobuchika Okada (KEK)**

# The cosmological constant problem

Observations of the Universe: **after WMAP results**

The flat universe dominated by **unknown energy densities**



**Dark energy: 73%**

**Dark Matter: 23%**

**Baryon: 4%**

## Standard 4D cosmology:

$$H^2 = \frac{1}{3\bar{M}_P^2} (\mathcal{E}_4 + \rho_D + \rho_B) \sim H_0^2 \sim (10^{-33} \text{eV})^2$$

Reduced Planck mass:  $\bar{M}_P \sim 10^{18}$  GeV

**Tiny, non-zero dark energy exist!**  $\mathcal{E}_4 \sim (10^{-3} \text{eV})^4$

If this is the cosmological constant  $\rightarrow$  **vacuum energy**

In field theories,  $\Delta V_0 \sim \Lambda_{cut}^4$

$\rightarrow$  **fine-tuning is necessary**

$$\mathcal{E}_4 = V_{tree} + \Delta V_0 = \mathcal{O}(\Lambda_{cut}^4) - \mathcal{O}(\Lambda_{cut}^4) \sim (10^{-3} \text{eV})^4$$

## The cosmological constant problem

- Why is it non-zero?
- Why is it so tiny  $\epsilon_4 \ll \Lambda_{cut}^4$  ?

Naturally,  $\Lambda_{cut} \sim M_P \sim 10^{19} \text{GeV}$

Introduction of **SUSY** may be a remarkable way,  
but it cannot solve the problem

If SUSY is manifest  $V_0 = \Delta V_0 = 0$

SUSY should be broken  $\rightarrow \epsilon_4 \sim M_{SUSY}^4 \geq (1 \text{TeV})^4$

It seems to be very hard to obtain a tiny vacuum energy

**Other possibilities?**

**Give up to reduce the vacuum energy,**

**but modify the 4D cosmological relation:**  $H^2 \sim \frac{\mathcal{E}_4}{\bar{M}_P^2}$

**4+N extra-dimensional theory:**

**the relation could be different**

**since  $\mathcal{E}_4$  could curve extra-dimensions**

**In this paper, the authors try to solve the cosmological constant problem based on the [brane induced gravity model with infinite extra-dimensions](#)**

# Modification of the cosmological relations

The author claim the modification

<b>In 4D</b>	$\rightarrow$	<b>In 4+N D</b>
$\bar{M}_P^2 H^2 \sim \mathcal{E}_4$		$M_*^{2+N} H^{2-N} \sim \mathcal{E}_4$

$M_*$  is the 4+N dim. Planck scale

**Example:**  $H \sim 10^{-33} \text{eV}$  ( $H^{-1} \sim 10^{28} \text{cm}$ )

$\rightarrow$

$N = 4, M_* \sim 10^{-3} \text{eV} \rightarrow \mathcal{E}_4 \sim (1 \text{TeV})^4$

$N = 6, M_* \sim 10^{-3} \text{eV} \rightarrow \mathcal{E}_4 \sim (\bar{M}_P)^4$

The model is very different from the usual extra-dim. models

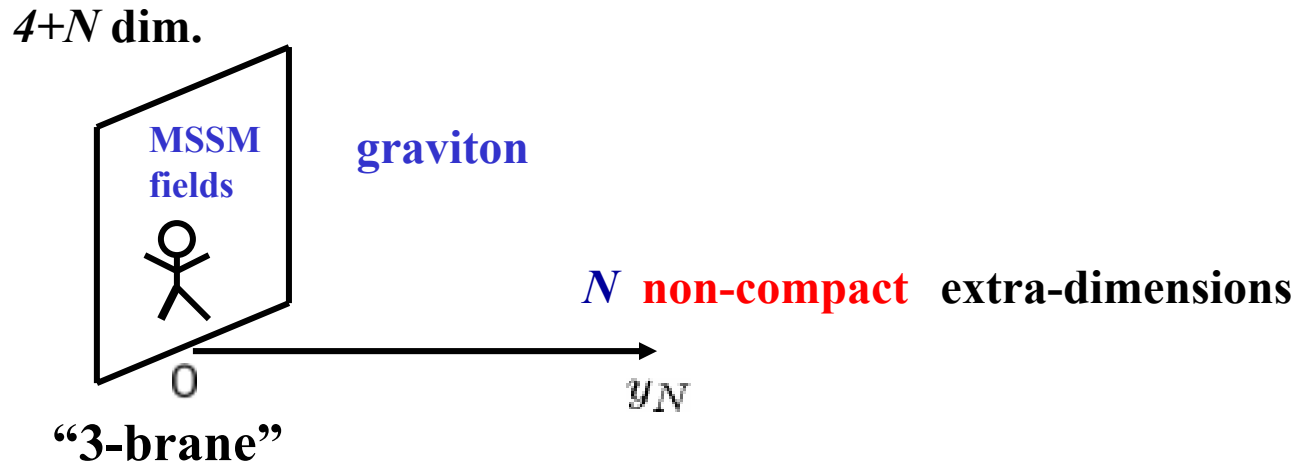
It is crucial that  $M_*$  is the model parameter independent of  $\bar{M}_P$

\*Remember the usual relation  $M_P^2 = M_*^{2+N} V_N$

# The model:

so-called ``brane induced gravity model''

Dvali, Gabadadze and Porrati,  
PLB485, 208 (2000)



$$S = M_*^{2+N} \int d^4x d^N y \sqrt{G} R + \int d^4x \sqrt{g} \left[ \Lambda_{cut}^2 \bar{R} + \mathcal{E}_4 + \mathcal{L}_{MSSM} \right]$$

I II III

# I: Bulk gravity with the 4+N dim. Planck scale independent of $\bar{M}_P$

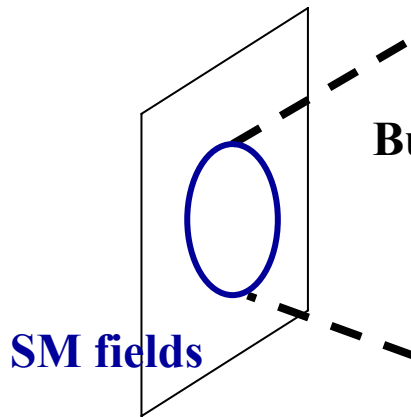
Bulk is supersymmetric  $\rightarrow \mathcal{E}_{bulk} = 0$

SUSY is broken on the brane  $\rightarrow$  but still  $\mathcal{E}_{bulk} = 0$

$$\Delta m_{SUSY} \propto V_N^{-1} \rightarrow 0$$

(infinite) volume suppression

## II: induced gravity on the brane



Bulk graviton

$$\sim \Lambda_{cut}^2 \bar{R}$$

is induced on the brane

We take  $\Lambda_{cut} \sim \bar{M}_P$

## III: cosmological constant on the brane

SUSY breaking on the brane  $\rightarrow \mathcal{E} \sim M_{SUSY}^4$



## Can the model be consistent with experiments?

Normally  $M_P^2 = M_*^{2+N} V_N \rightarrow \infty$  for infinite extra-dim.

→ inconsistent results

The brane induced gravity part  $S_{ind} = \int d^4x \sqrt{g} \bar{M}_P^2 \bar{R}$   
plays a crucial role

→ “Shielding effect”

for the observer on the brane, the gravity looks like

{ 4D gravity at short distance  
4+N dim. gravity at long distance

## Toy model: bulk scalar with kinetic term on the brane

$$S = M_*^{2+N} \int d^4x d^N y (\partial_M \Phi(x, y))^2 + \bar{M}_P^2 \int d^4x (\partial_\mu \Phi(x, y=0))^2$$

### Green's function

$$(M_*^{2+N} \partial_M \partial^M + \bar{M}_P^2 \delta^{(N)}(y) \partial_\mu \partial^\mu) G_R(x, y, 0, 0) = -\delta^{(4)}(x) \delta^{(N)}(y)$$

$$\rightarrow (M_*^{2+N} (p^2 - \Delta_N) + \bar{M}_P^2 p^2 \delta^{(N)}(y)) \tilde{G}_R(p, y) = \delta^{(N)}(y)$$

$$\tilde{G}_R(p, 0) = \frac{1}{\bar{M}_P^2 p^2 + M_*^{2+N} D^{-1}(p, 0)}$$

$$(p^2 - \Delta_N) D(p, y_N) = \delta^{(N)}(y)$$

$$D(p, y_N \rightarrow 0) \rightarrow 1/y^{N-2} \sim M_*^{N-2}$$

**Singularity at  $y=0$  would be softened by higher derivative terms for graviton ( $N>2$ )**

$$\tilde{G}_R(p, 0) \sim \frac{1}{\bar{M}_P^2 p^2 + M_*^4}$$

Therefore  $\tilde{G}_R(p, 0) \sim \frac{1}{\bar{M}_P^2 p^2}$  for  $p_c = 1/r_c \geq \frac{M_*^2}{\bar{M}_P}$

→ Green's function of 4D gravity is recovered

at a distance  $\leq r_c \sim \frac{\bar{M}_P}{M_*^2} \sim 10^{-33} \text{eV} \sim H_0^{-1}$

(for  $M_* \sim 10^{-3} \text{eV}$ )

Compare to usual model with  $\left\{ \begin{array}{l} \text{compactified extra-dim} \\ \text{no brane induced term} \end{array} \right.$

$$\left\{ \begin{array}{l} p > m_{KK} \rightarrow \text{N extra-dim. is going to be revealed} \\ p < m_{KK} \rightarrow \text{KK models are decoupled} \rightarrow \text{4D like} \end{array} \right.$$

## Cosmological solution with $\mathcal{E}_4 \neq 0$

Start with a static solution:

$$ds^2 = A^2(y)\eta_{\mu\nu}dx^\mu dx^\nu - B^2(y)dy^2 - C^2(y)y^2 d\Omega_{N-1}^2$$

$$A(y), B(y), C(y) \sim \left(1 - \left(\frac{y_g}{y}\right)^{N-2}\right)^\alpha \quad \text{was found}$$

Naked singularity appears at  $y = y_g \sim \frac{1}{M_*} \left(\frac{\mathcal{E}_4}{M_*^4}\right)^{\frac{1}{N-2}}$

because the solution is constrained  $\leftarrow$  static

→ The authors expect that the singularity would be removed if an inflation solution is considered

There are a number of examples, where a inflating solution removes the singularity with the relation

$$\frac{1}{H} \sim y_g$$

If we apply the relation to our case

$$H \sim M_* \left( \frac{M_*^4}{\mathcal{E}_4} \right)^{\frac{1}{N-2}}$$

**Example:**  $H \sim 10^{-33} \text{eV}$  ( $H^{-1} \sim 10^{28} \text{cm}$ )

$$\rightarrow \begin{aligned} N = 4, M_* \sim 10^{-3} \text{eV} &\rightarrow \mathcal{E}_4 \sim (1 \text{TeV})^4 \\ N = 6, M_* \sim 10^{-3} \text{eV} &\rightarrow \mathcal{E}_4 \sim (\bar{M}_P)^4 \end{aligned}$$

**Even if  $\mathcal{E}_4 \gg (10^{-3} \text{eV})^4$  we can obtain  $H \sim 10^{-33} \text{eV}$  consistent with the observations of the Universe**

## Summary

### Brane induced gravity model with infinite extra-dimensions

Three independent parameters:  $M_*$ ,  $\bar{M}_P$ ,  $\mathcal{E}_4$

4D gravity is obtained at a distance **shorter than**  $r_c \sim \frac{\bar{M}_P}{M_*^2} \sim H_0^{-1}$

Hubble parameter is obtained through  $H \sim M_* \left( \frac{M_*^4}{\mathcal{E}_4} \right)^{\frac{1}{N-2}}$