

Lectures on Brane World Scenarios II

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Large Extra Dimension Scenario

Phenomenology

- $M_{4+} \sim \mathcal{O}(1 \text{ TeV})$
- many graviton Kaluza-Klein modes

$$G_{\mu\nu}(x^\mu, y^1, y^2, \dots, y^n) = \sum_n g_{\mu\nu}^{(\vec{n})}(x^\mu) \chi^{(\vec{n})}(\vec{y})$$

$$\chi^{(n)} \propto e^{i\frac{\vec{n}\cdot\vec{y}}{r}}, \quad \left(m_{KK}^{(\vec{n})}\right)^2 = \frac{|\vec{n}|^2}{r^2} \quad \text{If 6 dim.} \rightarrow \frac{1}{r} \sim 10^{-4} \text{ eV}$$

Phenomenology of Extra-dimension Scenario

= Phenomenology of graviton Kaluza-Klein modes

Construction of effective action in 4D

1. 4D Lagrangian

4+ dim. Graviton \rightarrow 4D graviton in 4D

KK gravitons

KK gravi-scalars

KK gravi-vectors

2. Feynman rules

Characteristic features: infinite tower of KK gravitons

universal couplings

4D reduction of 4+ dim. Einstein's Eqs.

$$\square_{4+\delta} G_{MN} = -\frac{T_{MN}}{(M_{4+\delta})^{2+\delta}}$$

$$G_{MN} = \begin{bmatrix} G_{\mu\nu}^{(n)} & V_{\mu j}^{(n)} \\ V_{\mu i}^{(n)} & H^{(n)}\delta_{ij} + S_{ij}^{(n)} \end{bmatrix}$$

KK graviton $(\square + \hat{n})G_{\mu\nu}^{(n)} = \frac{1}{\bar{M}_P} \left[-T_{\mu\nu} + \frac{1}{3} \left(\frac{\partial_\mu \partial_\nu}{\hat{n}^2} \right) T_\lambda^\lambda \right]$

KK gravi-scalar $(\square + \hat{n})H^{(n)} = \frac{\sqrt{\frac{3(\delta-1)}{\delta+2}}}{\bar{M}_P} T_\mu^\mu$

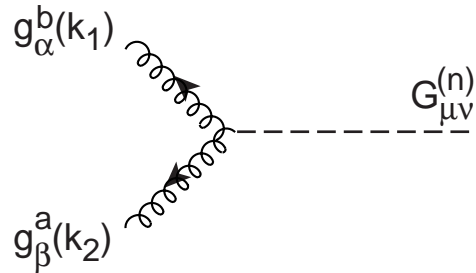
Vector $(\square + \hat{n})V_{\mu j}^{(n)} = 0$

Scalar $(\square + \hat{n})S_{ij}^{(n)} = 0$

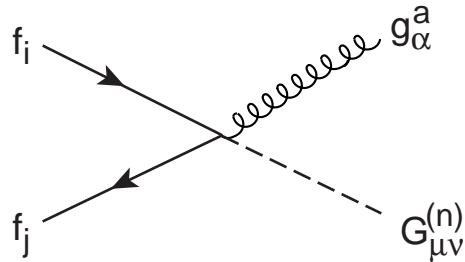
Feynman rules

$$T^{\mu\nu} = 2 \frac{\partial \mathcal{L}}{\partial g_{\mu\nu}} - g^{\mu\nu} \mathcal{L}$$

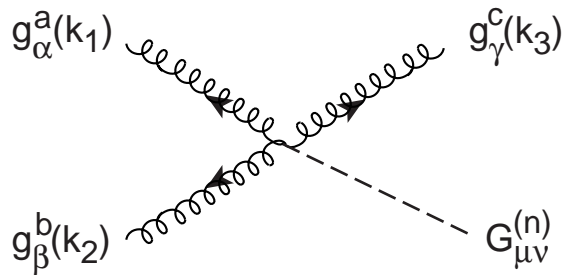
ex) QCD + KK graviton



$$-\frac{1}{\bar{M}_P} \left[\frac{i}{2} \eta_{\mu\nu} (k_{1\beta} k_{2\alpha} - (k_1 k_2) \eta_{\alpha\beta}) + \eta_{\alpha\beta} k_{1\mu} k_{2\nu} + \eta_{\mu\alpha} (\eta_{\nu\beta} (k_1 k_2) - k_{1\beta} k_{2\nu}) - \eta_{\mu\beta} k_{1\nu} k_{2\alpha} \right]$$



$$-\frac{i}{2\bar{M}_P} g_s t_{ji}^a [\gamma_\mu \eta_{\nu\alpha} + (\mu \leftrightarrow \nu)]$$



$$\frac{g_s}{\bar{M}_P} f^{abc} \left[Y(k_1)_{\mu\nu\alpha\beta\gamma} + Y(k_2)_{\mu\nu\beta\gamma\alpha} + Y(k_3)_{\mu\nu\gamma\alpha\beta} + Y(k_1)_{\nu\mu\alpha\beta\gamma} + Y(k_2)_{\nu\mu\beta\gamma\alpha} + Y(k_3)_{\nu\mu\gamma\alpha\beta} \right]$$

$$Y(k)_{\mu\nu\alpha\beta\gamma} = k_\mu (\eta_{\nu\beta} \eta_{\alpha\gamma} - \eta_{\nu\gamma} \eta_{\alpha\beta}) + [k_\beta (\eta_{\mu\alpha} \eta_{\nu\gamma} - \frac{1}{2} \eta_{\mu\nu} \eta_{\alpha\gamma}) + (\beta \leftrightarrow \gamma)]$$

Detection of Extra-dimension @ colliders

→ **detection of KK graviton**

direct → **KK graviton emission**

indirect → **KK graviton mediated process**

First detection of spin 2 particle !

KK graviton emission process @ hadron colliders

Emitted KK graviton \rightarrow non-interacting & stable particle
 \rightarrow **missing energy event**

For emissions of KK graviton with mass (m)

$$\frac{d\sigma_m}{dt}(q\bar{q} \rightarrow gG) = \frac{\alpha_s}{36} \frac{1}{s\bar{M}_P^2} F_1(t/s, m^2/s)$$

$$\frac{d\sigma_m}{dt}(qg \rightarrow qG) = \frac{\alpha_s}{96} \frac{1}{s\bar{M}_P^2} F_2(t/s, m^2/s)$$

$$\frac{d\sigma_m}{dt}(gg \rightarrow gG) = \frac{3\alpha_s}{16} \frac{1}{s\bar{M}_P^2} F_3(t/s, m^2/s)$$

$$F_1(x, y) = \frac{1}{x(y-1-x)} [-4x(1+x)(1+2x+2x^2) + y(1+6x+18x^2+16x^3) - 6y^2x(1+2x) + y^3(1+4x)]$$

$$F_2(x, y) = -(y-1-x)F_1\left(\frac{x}{y-1-x}, \frac{y}{y-1-x}\right)$$

$$F_3(x, y) = \frac{1}{x(y-1-x)} [1+2x+3x^2+2x^3+x^4 - 2y(1+x^3) + 3y^2(1+x^2) - 2y^3(1+x) + y^4]$$

$$\sum_n \rightarrow \int dn \rightarrow R^\delta \int d^\delta m$$

$$\frac{d\sigma}{dt} = S_{\delta-1} \left(\frac{\bar{M}_P^2}{M_{4+\delta}^{2+\delta}} \right) \int_0^{\sqrt{s}} m^{\delta-1} dm \times \left(\frac{d\sigma_m}{dt} \right)$$

$$S_{\delta-1} = \frac{2\pi^{\delta/2}}{\Gamma(\delta/2)} : \text{surface of a unit-radius sphere in } \delta \text{ dimensions}$$

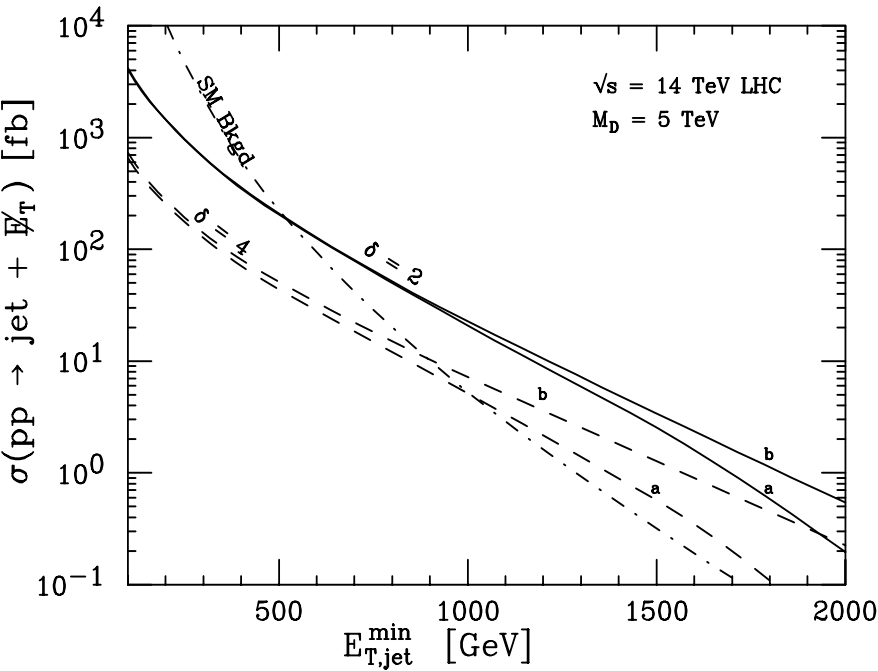
4D Planck scale is cancelled out

Cross sections depend on $\sigma \propto \left(\frac{\sqrt{s}}{M_{4+\delta}} \right)^{2+\delta}$

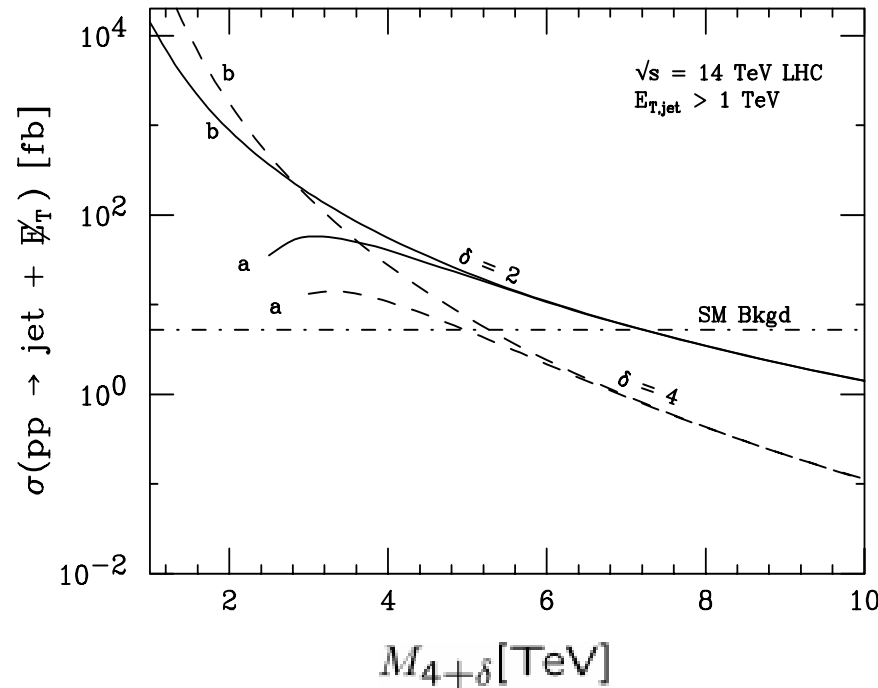
pp \rightarrow jet + nothing @ LHC $\sqrt{s} = 14\text{TeV}$

SM background = jet + Z ; Z \rightarrow two neutrinos

$M_{4+\delta} = 5\text{TeV}$



$E_{T,\text{jet}} > 1\text{TeV}$



a: integrating the cross section over $\hat{s} < M_{4+\delta}^2$

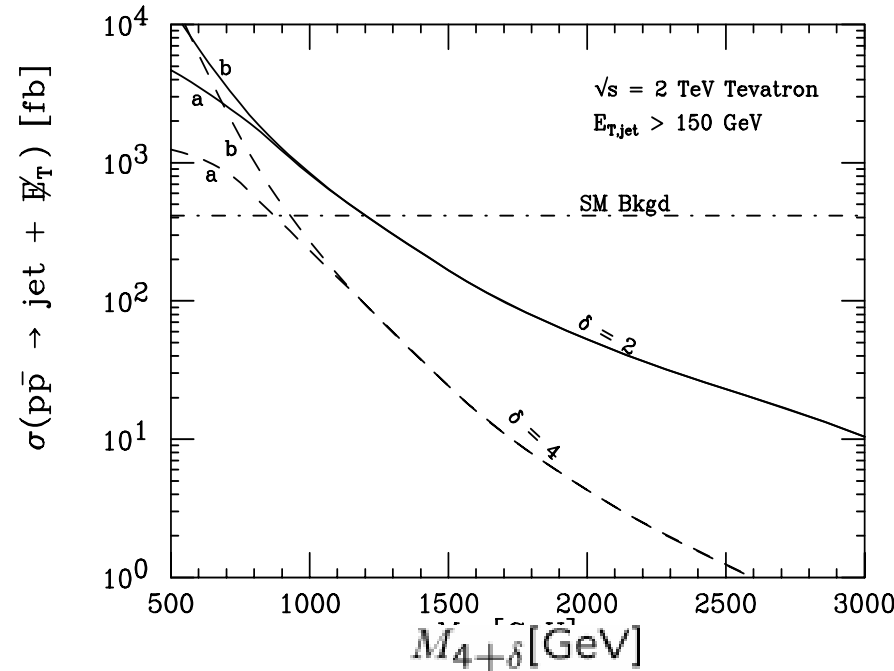
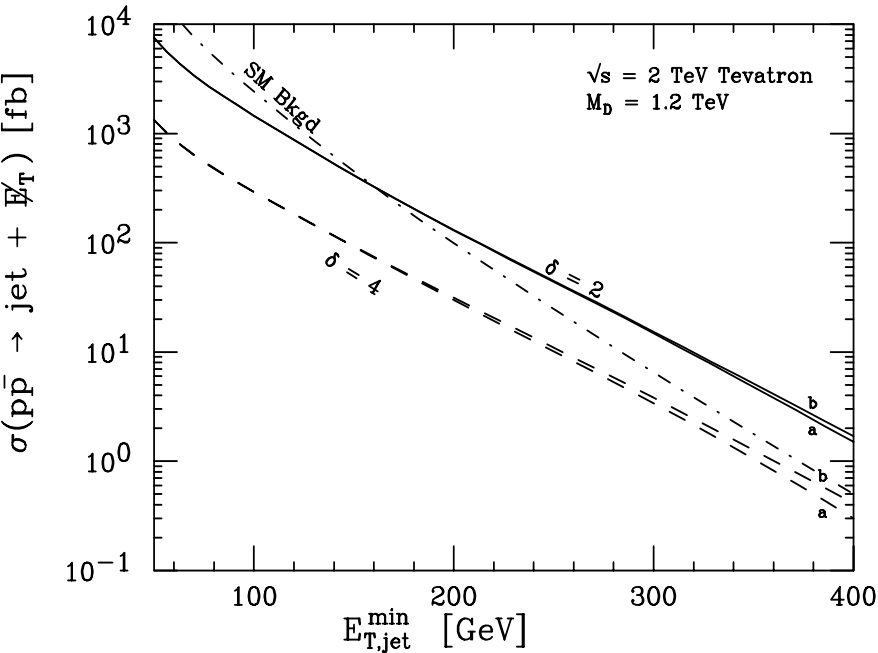
b: all

p pbar → jet + nothing @ Tevatron $\sqrt{s} = 2\text{TeV}$

SM background = jet + Z ; Z → two neutrinos

$$M_{4+\delta} = 1.2\text{TeV}$$

$$E_{T,\text{jet}} > 150\text{GeV}$$



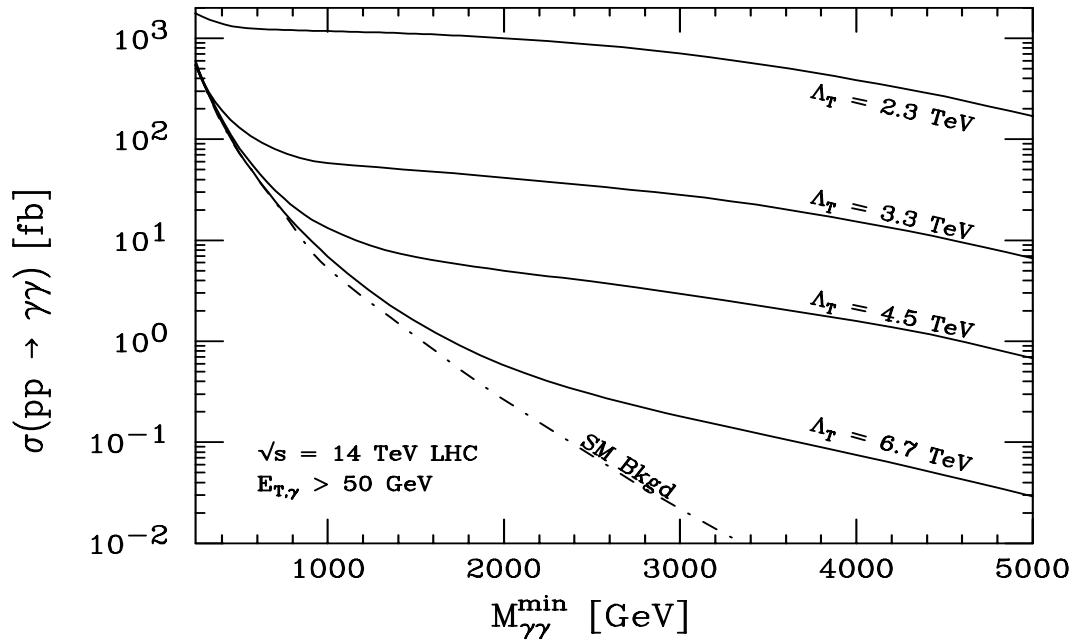
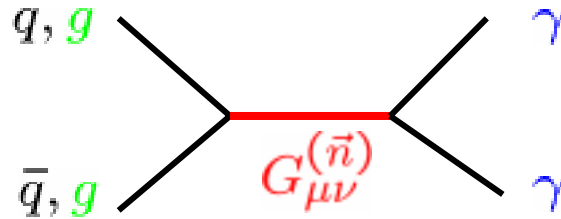
a: integrating the cross section over $\hat{s} < M_{4+\delta}^2$

b: all

Giudice et al., NPB544 (1999) 3

KK graviton mediated processes @ LHC

EX) $pp \rightarrow$



$$\sqrt{s} = 14 \text{ TeV}$$

$$E_{T,\gamma} > 50 \text{ GeV}$$

Giudice et al., NPB544 (1999) 3

$$\mathcal{M} = \frac{4\pi}{\Lambda_T^4} T_{\mu\nu}(p_1, p_2) T^{\mu\nu}(p_3, p_4)$$

Constraints from collider experiments

$\left\{ \begin{array}{l} \text{Lower bounds on } M_{4+} \\ \text{Upper bounds on } R \end{array} \right\}$
 from missing-energy exp.

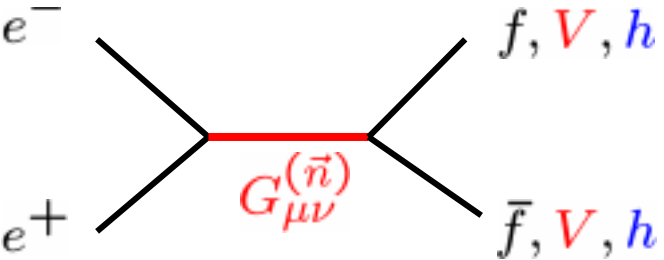
Example) LEP II $e^+e^- \rightarrow \gamma + \text{nothing}$
Tevatron $p\bar{p} \rightarrow \gamma + \text{nothing}$

Mirabelli, Perelstein & Peskin,
 Phys. Rev. Lett. 82 (1999) 2236

	$R(\text{cm})/M_{4+\delta}(\text{GeV})$		
LEP II	$4.8 \times 10^{-2}/1200$	$1.9 \times 10^{-9}/730$	$6.8 \times 10^{-12}/530$
Tevatron	$5.5 \times 10^{-2}/1140$	$1.4 \times 10^{-9}/860$	$4.1 \times 10^{-12}/780$
	$\delta = 2$	$\delta = 4$	$\delta = 6$

Updates by Cullen et al.,
 Phys. Rev. D62 (2000) 055012

Virtual KK graviton mediated processes in a soft brane scenario



$$\sum_{\vec{n}} \frac{1}{s - \left(m_{KK}^{(\vec{n})}\right)^2} \rightarrow \infty \quad (\text{for } \delta \geq 2)$$

Need regularization!

Naïve: Cut Off

$$m_{KK}^{MAX} \sim M_{4+\delta}$$

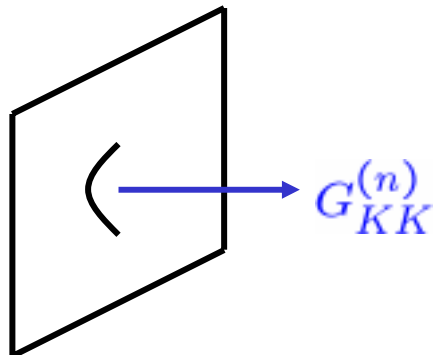
$$-\frac{8\pi}{M_P^2} \sum_{\vec{n}} \frac{1}{s - \left(m_{KK}^{(\vec{n})}\right)^2} \rightarrow \frac{4\pi\lambda}{M_S^4} \quad \lambda = \pm 1$$

More physically justified regularization?

Introduction of finite brane tension (= brane with finite width)

Brane fluctuation

Bando et al., PRL 83 (1999) 3601



Recoil effect

hard breaking of momentum conservation

Restoration of momentum

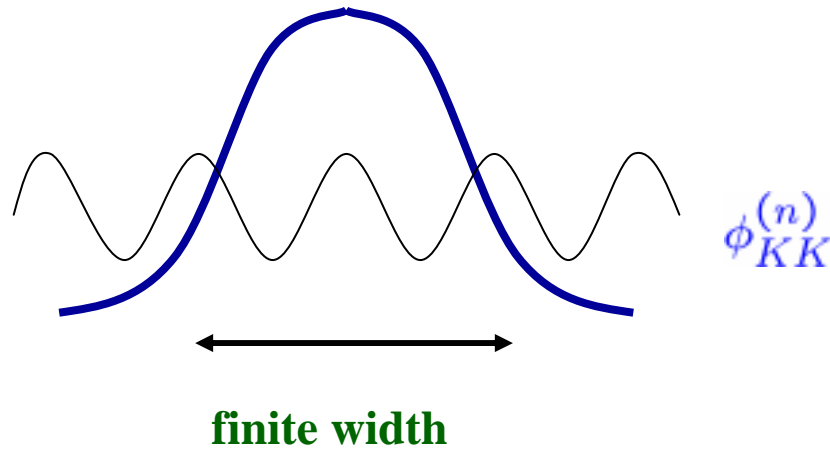
→ KK mode universal coupling

conservation for higher KK modes

$$g \longrightarrow g \times e^{-\frac{m_{KK}^2}{2\Delta^2}}$$

~ brane tension

More intuitive picture: brane with finite width



Hisano & N. Okada,
Phys. Rev. D61 (2000) 106003

Toy model:

brane fermion + bulk scalar

$$\psi \quad \phi_{KK}^{(n)}$$

Simple example: effective Yukawa coupling

$$Y_{eff} = Y_{univ} \int dy \psi(y) \phi_{KK}^{(n)} \psi(y)$$

$$\psi(y) \sim e^{-y^2/w^2}$$

$$\phi_{kk} \sim \cos(m_{KK}y)$$

$$\rightarrow Y_{eff} \sim Y_{univ} \times e^{-m_{KK}^2 w^2}$$

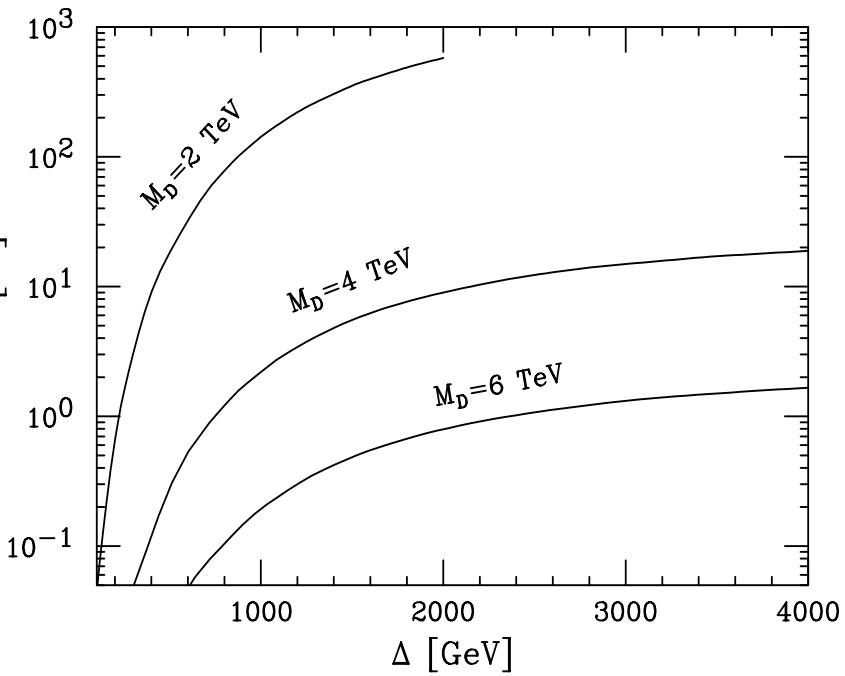
Restoration of momentum

conservation for higher KK modes

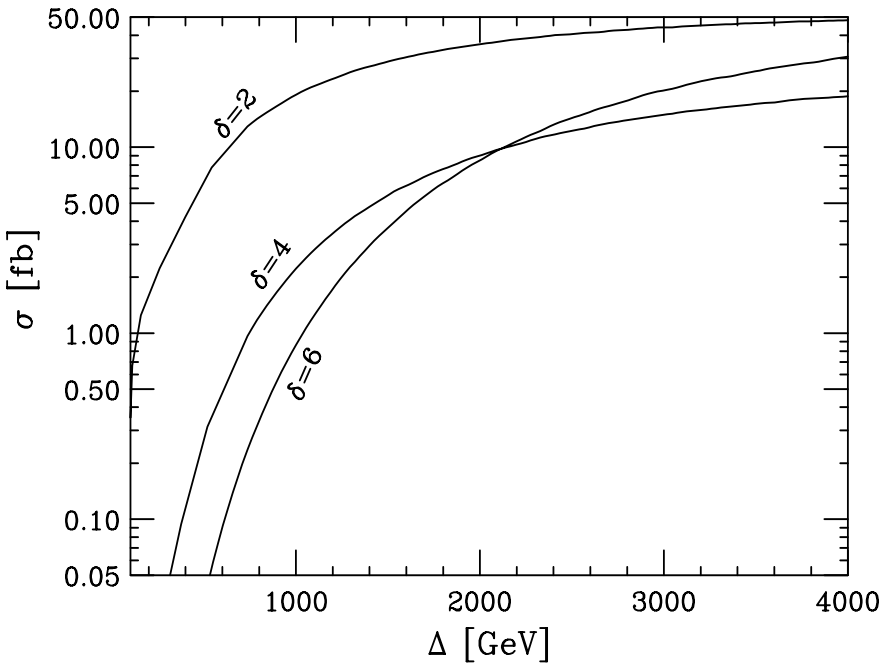
pp → jet + missing @ LHC

$$\mathcal{L} \sim -\frac{1}{\bar{M}_P} G_{\mu\nu}^{(n)} T^{\mu\nu} \times e^{-\frac{m^2(n)}{2\Delta^2}}$$

$\delta = 4$



$E_{T,jet} > 1\text{TeV}$

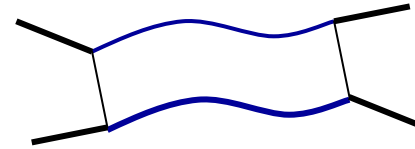
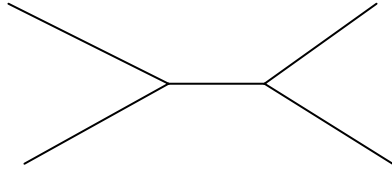


$M_{4+\delta} = 4\text{TeV}$

TeV scale string theory

Evidence of TeV scale string theory?

$$M_{4+n} \sim M_{st} \sim 1 \text{ TeV}$$



String amplitude

(Perturbative) String Theory

Many states: **string Regge excitations**

$$M_n = \sqrt{n} M_{st} \quad (n = 0, 1, 2, \dots)$$
$$J = 0, 1/2, 1, 3/2, 2, 5/2, \dots$$

Phenomenology of TeV string theory

string Regge excitations: **new contact terms**
string resonances

$$\sqrt{s} \leq M_{st}$$

$$\sqrt{s} \geq M_{st}$$

Stringy correction to scattering amplitudes

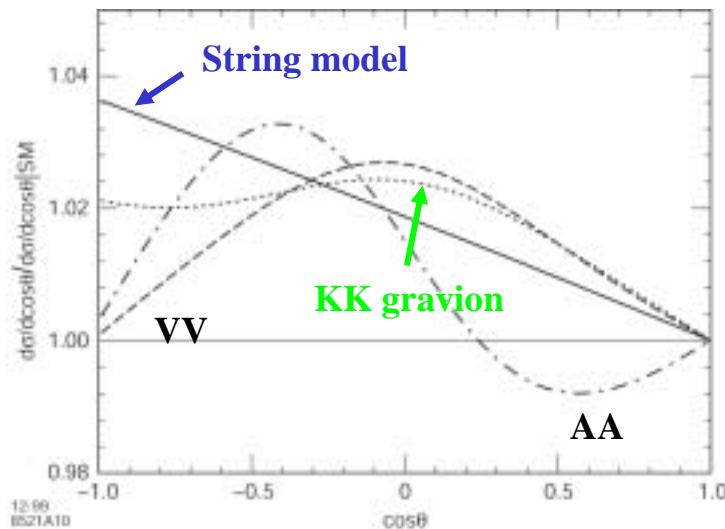
Cullen, Perelstein & Peskin,
PRD 62 (2000) 0550212

$$\frac{d\sigma}{d\cos\theta} = \left(\frac{d\sigma}{d\cos\theta}\right)_{SM} |S(s, t)|^2$$

String form factor: $S(s, t) = \frac{\Gamma(1 - s/M_{st}^2)\Gamma(1 - t/M_{st}^2)}{\Gamma(1 - s/M_{st}^2 - t/M_{st}^2)}$

$\sqrt{s} \leq M_{st} \rightarrow$ Contact interaction mediated SR excitations

Example: Bhabha scattering @ 1TeV LC



$$M_{st} = 3.1\text{TeV}$$

$$\Lambda_T = 6.2\text{TeV}$$

VVcontact with $\Lambda = 88\text{TeV}$

AAcontact with $\Lambda = 62\text{TeV}$

$$\mathcal{L}_{KK} = i\frac{4}{\Lambda_T} T^{\mu\nu} T_{\mu\nu}$$

$$\mathcal{L}_{VV} = -\frac{4\pi}{2\Lambda^2} \bar{e}_L \gamma^\mu e_L \bar{e}_L \gamma_\mu e_L + \dots$$

$s = M_n^2 = nM_{st}^2 \rightarrow$ String resonance

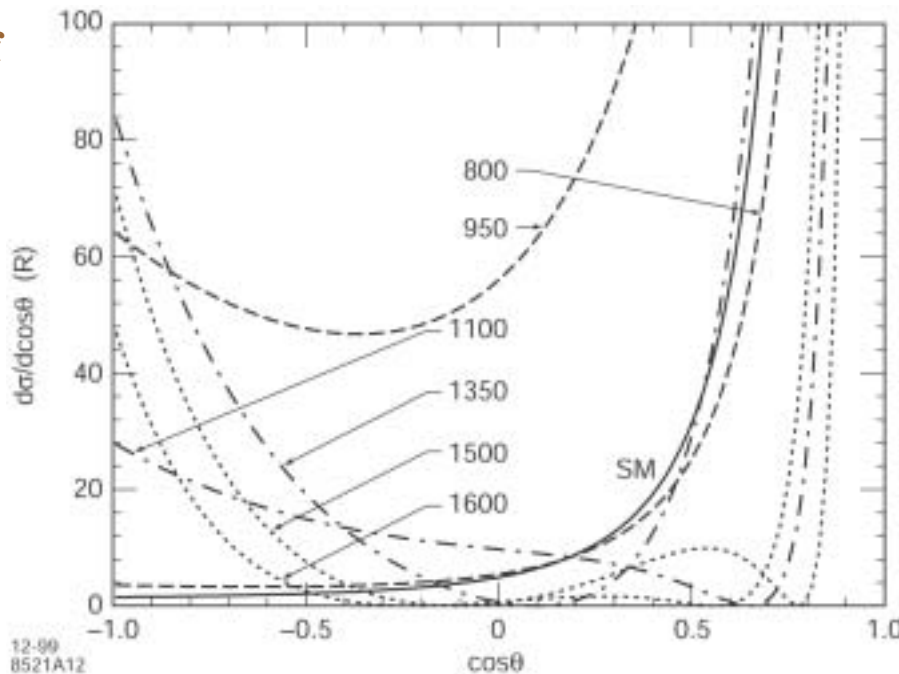
$nM_{st}^2 < s < (n + 1)M_{st}^2$ behavior of string form factor

Bhabha scattering @ 1TeV LC

Cullen, Perelstein & Peskin,
PRD 62 (2000) 0550212

In units of

$$\frac{4\pi\alpha^2}{3s}$$



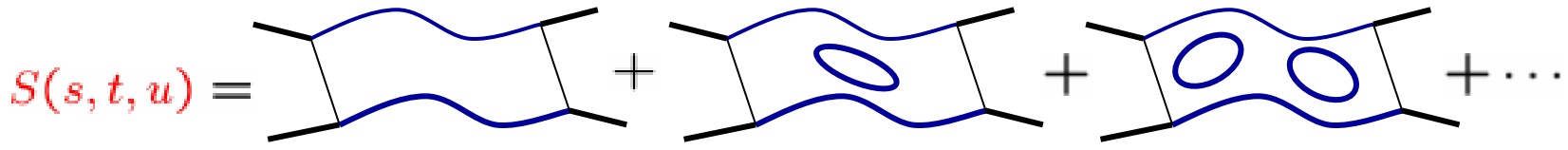
$$M_{st} = 1\text{TeV}$$

$$s < M_{st}^2$$

$$M_{st}^2 < s < 2M_{st}^2$$

$$2M_{st}^2 < s < 3M_{st}^2$$

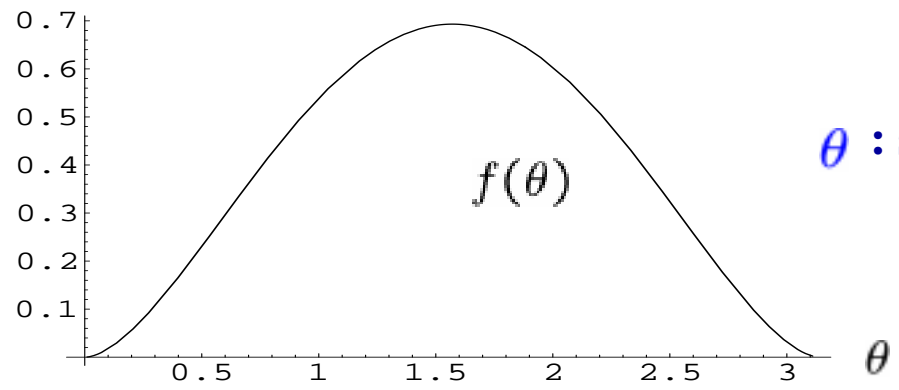
Stringy effect: hard scattering limit



$$\sim e^{-\sqrt{\frac{sf(\theta)}{M_{st}^2}}}$$

$$f(\theta) = -\lambda \ln(\lambda) - (1 - \lambda) \ln(1 - \lambda)$$

$$\lambda = \sin^2(\theta/2)$$



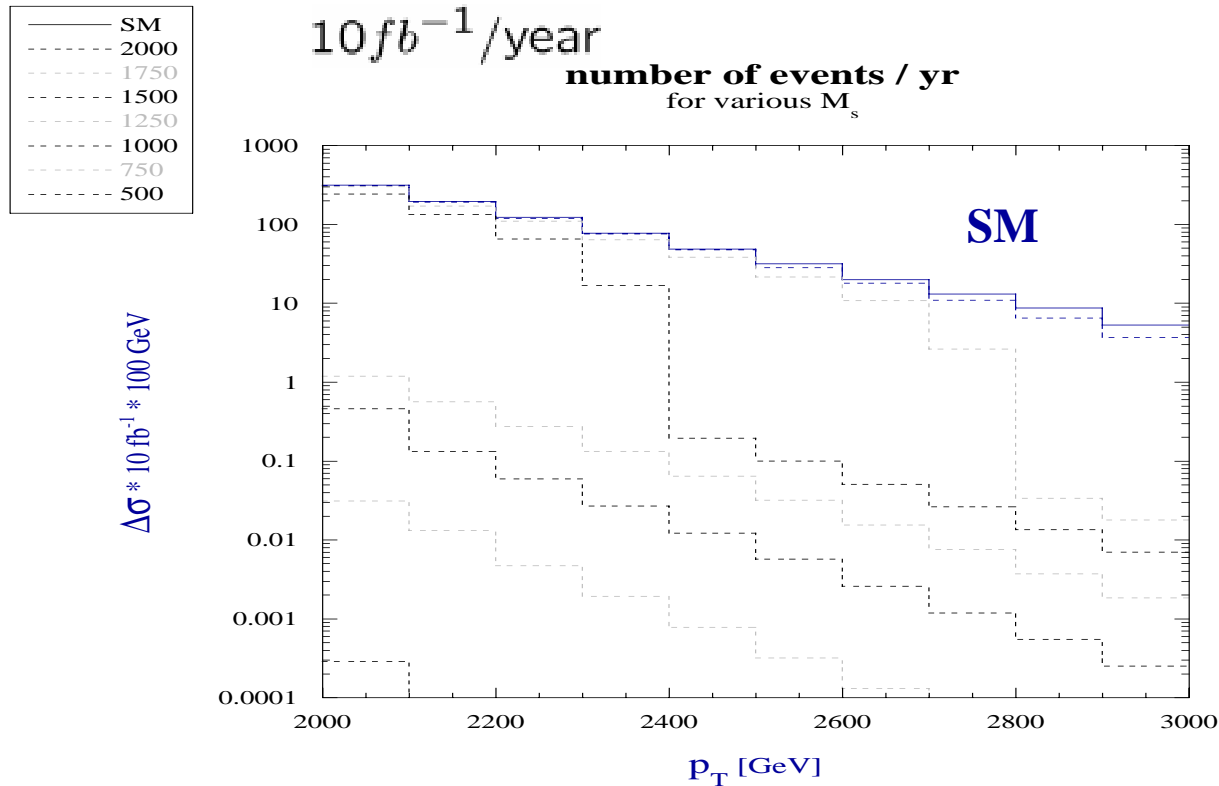
θ : scattering angle

Stringy behavior @ LHC, if $M_{st} = O(1 \text{ TeV})$

$$M_{st} \ll \sqrt{s_{LHC}}$$

Large angle scattering $\rightarrow \sigma \sim \sigma_{SM} \times e^{-\frac{s}{M_{st}^2}}$

Dramatic reduction in QCD jet production @ LHC



K. Oda & N. Okada,
PRD 66 (2002) 095005

Non-perturbative process?

$$\sqrt{s} > M_{4+\delta}$$

TeV string, strong gravity

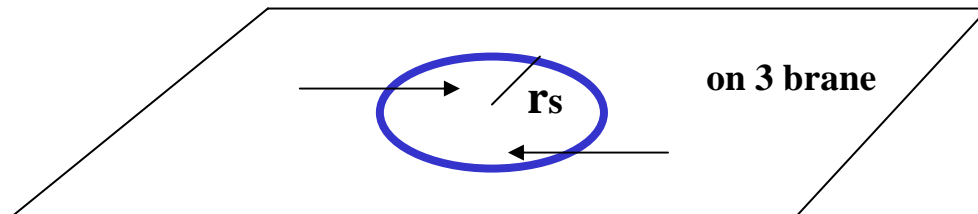
4+ dim. Black Hole formation?

Giddings-Thornas, PRD65 (2002) 056010
Dimopoulos-Lansberg, PRL87 (2001) 161602

4+ dim. BH solution

$$ds^2 = -R(r)dt^2 + R(r)^{-1}dr^2 + r^2 d\Omega_{\delta+2}^2$$

$$\begin{cases} R(r) = 1 - \left(\frac{r_s}{r}\right)^{\delta+1} \\ r_s = \frac{1}{\sqrt{\pi}M_{4+\delta}} \left(\frac{8\Gamma\left(\frac{\delta+2}{2}\right)}{\delta+2}\right)^{\frac{1}{\delta+1}} \left(\frac{M_{BH}}{M_{4+\delta}}\right)^{\frac{1}{\delta+1}} \end{cases}$$



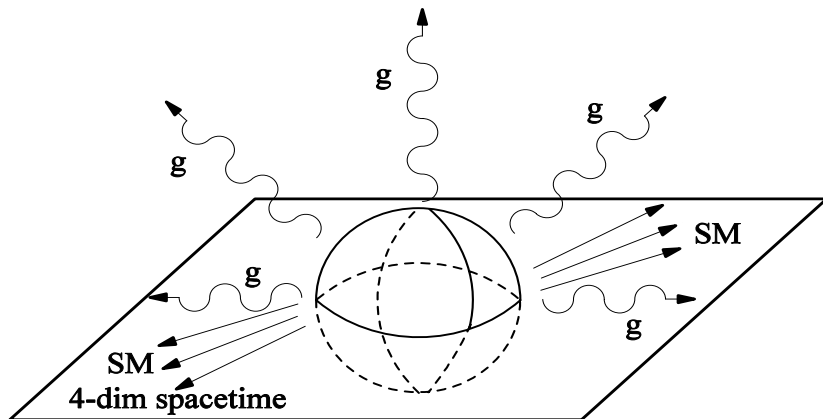
$$\sigma(M_{BH}) \sim \pi r_s^2$$

BH production @ LHC

$$\sigma_{pp \rightarrow BH}(x_m, \delta) \sim \sum_{ij} \int_{x_m}^1 dx \int_x^1 \frac{dy}{y} f_i(y, Q) f_j(x/y, Q) \times \sigma_{ij \rightarrow BH}(xs; \delta)$$

$$x_m = \frac{M_{BH}^{min}}{\sqrt{s}} \sim \frac{M_{4+\delta}}{\sqrt{s}}$$

BH decay → Evaporate by emitting Hawking Radiation



$$T_H = \frac{\delta + 1}{4\pi r_s}$$

- particle blind
- $E \sim T_H$ for each quanta
- emission rate → # of d.o.f

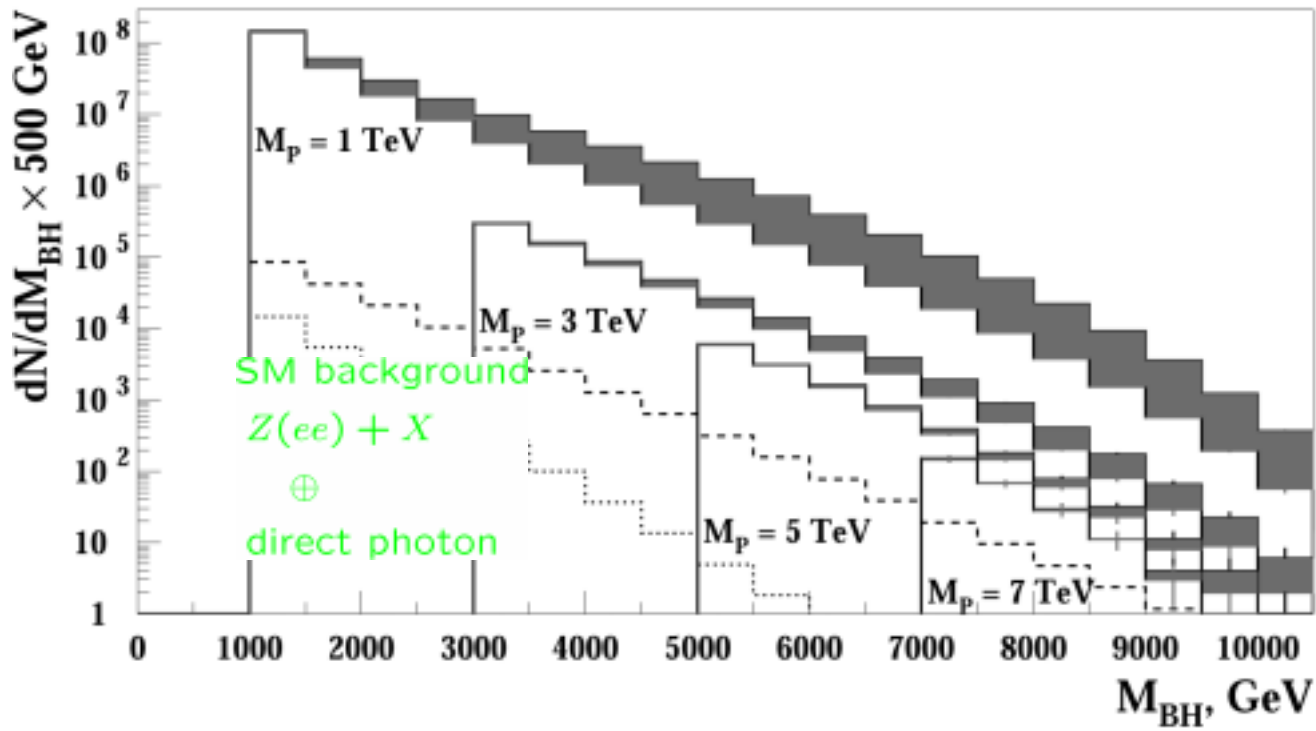
BH →

- quarks & gluons ~ 75 %
- charged leptons ~ 10 %
- neutrinos ~ 5 %
- , W, Z ~ 2 % each

of produced BH @ LHC ($100 fb^{-1}$ integrated luminosity)

electron/photon decay channels

Dimopoulos-Landsberg,
PRL 87 (2001) 161602



$\delta = 1 \sim 7$

LHC can be “BH Factory”

if $M_{BH}^{min} \sim \mathcal{O}(1\text{TeV})$

Far future

BH property has been known

Collider will be designed in order to

NOT produce BH