Beyond the Standard Model: the next 20 years

1. (Orientation and) LHC
Many people interested in experimental elementary particle physics are interested in many different problems, including flavor, neutrinos, CP violation, proton decay, and grand unification, and aspects of inflationary cosmology.

With my apologies, I am only interested in one problem -- the study of new particles and forces in the energy range 100 GeV - 1 TeV.

I am convinced that these new particles exist. If this is so, we will begin to discover them in the next few years.

And, if this is so, this study will define elementary particle physics for the next 20 years.
In these lectures, I would like to discuss three chapters in this study:

1. **LHC**: How will we discover the existence of new particles in the hundred-GeV mass range?

2. **ILC**: How will we study the interactions of these particles in detail, and what will we learn by doing this?

3. **Beyond ILC**: How will we pursue the study of these particles to higher energies?
The presence of new physics in the hundred-GeV mass range is motivated by two pressing problems:

electroweak symmetry breaking and cosmic dark matter

I will review the first of these today, the second tomorrow.
The problem of electroweak symmetry breaking is particularly compelling.

We know that the weak and electromagnetic interactions are based on an $SU(2) \times U(1)$ Yang-Mills gauge theory.

This theory prohibits masses for the quarks, leptons, $W$ and $Z$ bosons unless the symmetry is spontaneously broken.

There is a simple theory of this spontaneously symmetry breaking based on the minimal Higgs boson. But this theory is inadequate. It is not a physics explanation.

Let me review the pieces of this story.
First, the SU(2) x U(1) structure of the weak interactions is well established experimentally. This structure is built on three key observations:

1. **Universality**: all QED, W, Z couplings arise from the two constants $g, g'$

2. **Chirality**: left- and right-handed fermions have different weak couplings

3. **non-Abelian gauge structure**: The couplings of vector particles are of the form predicted by Yang and Mills.

All of these features received new confirmation by the LEP and SLC experiments.
Universality:

In the context of the precision Z experiments, this is the statement that the coupling of each species of quark and lepton is given by

$$\frac{g}{c_w}(I^3 - s^2_w Q)$$

with only two parameters $g, s^2_w$. These parameters can be extracted from $\alpha, G_F, m_Z$, e.g.,

$$\sin^2 2\theta_w = \frac{4\pi\alpha(m_Z^2)}{\sqrt{2}G_F m_Z^2}$$

These points are tested by the measurements of the Z partial widths.
The partial width of the Z into a fermion species $f$ should be given by:

$$\Gamma(Z \rightarrow f\bar{f}) = \frac{\alpha m_Z}{6s_w^2 c_w^2} \cdot S_f$$

times the factor $3(1 + \alpha_s/\pi) = 3.11$ for quarks.

This gives the following table of partial widths and branching ratios:

<table>
<thead>
<tr>
<th>species</th>
<th>$\Gamma(Z \rightarrow f\bar{f})$</th>
<th>BR</th>
</tr>
</thead>
<tbody>
<tr>
<td>$\nu_e, \nu_\mu, \nu_\tau$</td>
<td>167 MeV</td>
<td>6.7%</td>
</tr>
<tr>
<td>$e, \mu, \tau$</td>
<td>84 MeV</td>
<td>3.4%</td>
</tr>
<tr>
<td>$u, c$</td>
<td>300 MeV</td>
<td>12.0%</td>
</tr>
<tr>
<td>$d, s, b$</td>
<td>383 MeV</td>
<td>15.3%</td>
</tr>
</tbody>
</table>

Including a small correction for the case of $\Gamma(Z \rightarrow b\bar{b})$, we find a total width

$$\Gamma_Z = 2.50 \text{ GeV}$$
To test these predictions, we first measure e+e- annihilation at the Z resonance and measure the relative branching ratios to hadrons and to visible leptons.

Then we must determine the total width.

The shape of the resonance is distorted by initial-state photon radiation. Thus, it is necessary to measure the detailed shape of the resonance to extract $\Gamma_Z$.

It is amusing to note that all three of the Standard Model interactions - QED, QCD, and of course SU(2) x U(1) - contribute to the Z line-shape.

The result is: $\Gamma_Z = 2.4952 \pm 0.0023$ MeV
There is a special consideration for the $b$ quark. The diagrams contribute a correction to the $b_L$ $Z$ charge,

$$Q_{ZbL} = -\left(\frac{1}{2} - \frac{1}{3} s_w^2 - \frac{\alpha}{16\pi s_w^2} \frac{m_t^2}{m_W^2}\right)$$

This is a -2% correction to the partial width. It is easier to measure the quantity

$$R_b = \frac{\Gamma(Z \rightarrow b\bar{b})}{\Gamma(Z \rightarrow \text{hadrons})}$$

which, if universality is correct, is almost independent of $s_w^2$. 
The final result is:

\[ R_b = 0.21643 \pm 0.00073 \]

in excellent agreement with the Standard Model and confirming the -2\% shift due to the t-W diagrams.
Chirality:

It is clear from phenomenology of parity violation in beta decay that the weak interactions couple differently to left- and right-handed fermions.

In the context of a gauge theory, this tells us directly that the left and right species have different gauge quantum numbers.

There is a new test of this in the Z decays, the measurement of final-state helicity in $Z^0 \rightarrow f \bar{f}$. This is given by

$$A_f = \frac{g_L^2 - g_R^2}{g_L^2 + g_R^2}$$

$A_f$ has quite different values for different species:

- 15% for leptons,
- 94% for down-type quarks
From this and other measurements of final-state lepton polarization, we obtained:

\[ A_\ell = 0.1465 \pm 0.0033 \]

It was also possible at SLAC to polarize the electrons and measure \( A_e \) directly as an asymmetry in the total cross section on the Z resonances. This gave:

\[ A_e = 0.1513 \pm 0.0021 \]
$A_b = 0.94$ at the $Z^0$ SLD
Yang-Mills structure:

This is tested by the LEP measurements of $\sigma(e^+e^- \rightarrow W^+W^-)$. This reaction has a long-recognized danger of violation of unitarity.

$W^+$ has 3 polarization states. In the rest frame $\epsilon^\mu = (0, \hat{n})^\mu$

but for a $W$ in motion

$$p^\mu = (E_W, 0, 0, k_W)$$

$$\epsilon_R = \frac{1}{\sqrt{2}}(0, 1, i, 0)$$

$$\epsilon_L = \frac{1}{\sqrt{2}}(0, 1, -i, 0)$$

$$\epsilon_0 = \frac{1}{m_W}(k_W, 0, 0, E_W) \approx p^\mu / m_W$$

Notice that

$$\epsilon^*_+ \cdot \epsilon^*_- = \frac{E_W^2 + k_W^2}{m_W^2} \approx \frac{s}{2m_W^2}$$

This is trouble; unitarity requires $|iM(e^+e^- \rightarrow W^+W^-)| < \text{const}$ in each partial wave.
However, in a spontaneously broken gauge theory, the 0 polarization state of the $W$ comes from eating a Goldstone boson. It turns out that the predicted cross section is just that for producing the Goldstone bosons.

$$i \mathcal{M}(e_L^− e_R^+ \rightarrow W^+_0 W^-_0) =$$

$$= ie^2 \left[ \frac{1}{4c_w^2} + \frac{1}{4s_w^2} \right] \bar{v} \gamma^\mu u \frac{1}{s} (k_+ - k_-)_\mu$$

The SU(2) x U(1) model gives this result by a delicate cancellation among the diagrams.

This cancellation takes place only if the form of the 3-boson vertex is exactly that given by Yang-Mills theory.
$\sigma(e^+ e^- \rightarrow W^+ W^-)$

**LEP**
PRELIMINARY

**ALEPH**

![Graph showing $\sigma(e^+ e^- \rightarrow W^+ W^-)$ as a function of Ecm (GeV)]

- **YFSWW/RacoonWW**
- **no ZWW vertex (Gentle)**
- **only $\nu_e$ exchange (Gentle)**

**c)**

![Histogram showing the distribution of $\cos(\theta^*)$]
Thus, the evidence that the weak and electromagnetic interactions are a

spontaneously broken gauge theory of SU(2) x U(1)

is impressively strong.

This brings up an obvious question:

What is the explanation for the breaking of SU(2) x U(1)?
In the minimal form of the Standard Model, we explain this spontaneous symmetry breaking by postulating the existence of an SU(2)-doublet scalar field $\phi$ with the potential

$$V = \mu^2 |\phi^2| + \lambda |\phi|^4$$

If $\mu^2 < 0$, the minimum of the potential does not respect $\text{SU}(2) \times \text{U}(1)$.

So, why is $\mu^2 < 0$? **No answer!**

$\mu^2$ receives additive corrections from higher-order corrections

$$\mu^2 = \mu^2_{\text{bare}} + \frac{\lambda}{8\pi^2} \Lambda^2 - \frac{3y_t^2}{8\pi^2} \Lambda^2 + \cdots$$

So $\mu^2 < 0$ is not a simple criterion in the underlying theory.
What does an explanatory theory of SU(2) x U(1) breaking look like?

The theory must include an SU(2) doublet that obtains a vacuum expectation value. This field can be either composite or (effectively) elementary at the 100 GeV scale.

If the doublet of fields is composite, the theory should include their excited states. It is very difficult for these not to upset the precision electroweak results. So I will assume that there is an elementary Higgs field $\phi$.

The theory should generate the potential for $\phi$ from physics. That is,

1. $\mu^2$ should not receive additive, divergent corrections.
2. A calculation should give $\mu^2 < 0$. 
It is very difficult to prohibit additive corrections to the mass term of a scalar field. However, there are three known ways to forbid this term by symmetry:

\[ \delta \phi = \epsilon \quad \phi \text{ is a Goldstone boson} \]

\[ \delta \phi = \epsilon \cdot A \quad \phi \text{ is part of a higher-dimensional gauge field} \]

\[ \delta \phi = \epsilon \cdot \psi \quad \phi \text{ is part of a supersymmetry multiplet} \]

In each case, there is a natural mechanism to generate a potential with \( \mu^2 < 0 \), if the top quark Yukawa coupling is the largest relevant coupling in the model.
In Little Higgs models, the loop corrections due to the top quark and its partner cancel with a negative residue.

In extra-dimensional models, the Kaluza-Klein excitations of the top quark give a symmetry-breaking potential for $A^5$

In supersymmetry, the renormalization by the top quark Yukawa coupling gives a negative correction.

In all cases, we need a complex model, with new particles that are partners of the top quark.
I do not ask you to literally accept these models, but I do ask you to accept the principle that a model of electroweak symmetry breaking must have multiple components and interacting parts.

These cannot belong to the Standard Model. They are new particles associated with the hundred-GeV mass scale of the Higgs potential.

By this logic, the new particles must be there.

If some are partners of the top quark, they must have QCD interactions. Can we find them?
arrival of a superconducting muon toroid at CERN

Paula Collins, CERN
CMS iron toroid installation
As I will discuss in a moment, the LHC brings with it an exceptionally difficult environment in which to search for new physics. But this does not mean that new physics must be hard to find. There are two scenarios that are quite likely in which the discovery is straightforward.

new vector bosons decaying to $\mu^+ \mu^-$

new stable heavy leptons

However, if you will excuse me, I will concentrate in this lecture on the more generic hypothesis of new heavy particles with QCD color. I will return to these models in lecture 3.
One feature that I would like to keep in this discussion is the possibility of producing invisible particles that carry away missing energy and momentum. I will argue tomorrow that the fundamental particle of cosmic dark matter is likely to be in the hundred-GeV mass range. This would be a final decay product of new particles produced at the LHC.

For new particles with QCD color, the expected cross sections are at the 10 pb level, corresponding to $10^5$ events/yr at the luminosities expected for early LHC running.
The new physics events should be quite complex. A typical event would have the form shown. Particle labels are for supersymmetry, but this type of event can appear in all three scenarios discussed earlier.

It is expected that events of this kind will appear as a very significant signal above background.

Here are the estimates of Tovey (2003) for supersymmetry models with universal scalar and gaugino masses at the GUT scale.
For squark and gluino masses below 1 TeV, the missing energy signature should be significant with a very small amount of integrated luminosity.
At the same time, many different signatures of new physics should be seen above the Standard Model expectation.
However, the expectation of large signals above Standard Model background does not mean that we can be complacent.

The theoretical background levels must be understood very well both absolutely and in relation to the actual data.
Gianotti and Mangano (2005):

“Not only is the rate larger than previously expected, but the shape of the distribution is different, and much closer to that of the signal itself.”
An enormous amount of work has been done on the theoretical calculation of these background rates.

But still all particle physicists - even string theorists - should be engaged with this problem. We need clearer ways to think about the prediction of backgrounds, and to verify our models of them from data.
In order to reach the level of new physics signals, we will need to work down through a series of levels dominated by Standard Model processes of different types.

Here is an idea of the hierarchy:

\[
\begin{align*}
\sigma_{tot} & \quad 100 \text{ mb} \\
\text{jets w. } p_T > 100 & \quad 1 \mu \text{b} \\
\text{Drell-Yan} & \quad 100 \text{ nb} \\
\bar{t}t & \quad 800 \text{ pb} \\
\text{SUSY } (M < 1 \text{ TeV}) & \quad 1-10 \text{ pb}
\end{align*}
\]
The first challenge comes with the realization that the processes that we are looking for occur at rates of order

$$10^{-11}$$
of the total pp cross section.

Still, the interesting events have several jets with large values of pT. To find jets, we can look at the ‘lego plot’ of pT deposited in the plane of $\theta$ and $\phi$ - or, better, rapidity $y$ and $\phi$. If we look for these objects instead of simply searching for large energy deposition, we already win about 6 orders of magnitude.
Run 223385 Evt 9802792 Thu Jul 20 17:14:11 2008

ET scale: 10 GeV

DO event
lego plot of DO event
ET scale: 436 GeV

DO event
lego plot of DO event

Bins: 481
Mean: 2.32
Rms: 23.9
Min: 0.00933
Max: 384
To go further, we need to search for events that do not belong to the classes generated by QCD. These should be events with multiple jets, plus leptons or unbalanced visible momentum.

QCD will generate unbalanced momentum if jets are mismeasured. To control this effect, it is necessary to understand the detectors, to eliminate noise and electronic signals unrelated to the physics events, and to correct for cracks and geometric inefficiencies.
CMS and ATLAS claim that they can control these effects to the required level. That story is expressed in these figures from the ATLAS TDR.

ATLAS simulation of missing ET in $Z(\rightarrow \mu^+\mu^-) + \text{jet}$

$\eta$ of the jet w. the highest ET in events w. ET > 50
In the physics studies of ATLAS and CMS, the dominant backgrounds to new physics come from a different source, heavy particle production within the Standard Model, production of $W, Z, t\bar{t}$ plus jets.

These reactions already offer missing energy, leptons, and hadronic activity. They populate the region of large HT associated with new physics to the extent the additional jets are radiated along with the heavy particles.
This is genuinely scary. Processes such as

\[ \nu \rightarrow \bar{\nu} \]

have cross sections comparable to the SUSY signal and might compete with it.
Here is a recent quantitative evaluation by Sanjay Padhi, using ALPGEN and the ATLAS full simulation code.

\[ E_T > 100 \]
\[ 4 \text{ jets, 2 w.} \]
To understand heavy particle + multijet backgrounds to new physics, there is a methodology that has been used successfully in the Tevatron, especially in the CDF and DO analyses of top quark production.

Use the fact that new particles appear in events with large numbers of jets and large

\[ H_T = \sum_i E_{T_i} \]

Compute systematically the SM rates for n jet production. The results for fewer jets can be validated against data, both in a general setting and also with the experimental cuts that define the new physics search. Now extrapolate to large numbers of jets and large \( H_T \).
This method is now a standard part of the Tevatron culture.

It apparently originated in UA2, where the systematics of jet counting was called “Berends scaling”. The name did not stick, and there are earlier references.

I think that the concept—in its original context, and in greater generality—is very important for carrying out and evaluating experimental results from the LHC. I would like to present a new name for it: the staircase.

Here is—to my knowledge—the original staircase presented by Ellis, Kleiss, and Stirling:
Ellis-Kleiss-Stirling staircase (1985)

compared to preliminary data from UA1.
Here is the published UA1 data, compared to a calculation based on the Berends-Giele technology for multijet computation:

UA1
1988

Berends, Giele, Kuijf, Kleiss, Stirling
1989
Let me show you a series of recent figures from the Tevatron experiments that illustrate this concept.
systematics of $W + \text{jets}$

$$(W \rightarrow e\nu) + \geq n \text{jets}$$

CDF Run II Preliminary

CDF Data $\int dL = 320 \text{ pb}^{-1}$

$W_{\text{kin}}$: $E_T^e \geq 20[\text{GeV}]$; $|\eta^e| \leq 1.1$

$M_T^W \geq 20[\text{GeV/}c^2]$; $E_T^\nu \geq 30[\text{GeV}]$

Jets: JetClu R=0.4; $|\eta|<2.0$

hadron level; no UE correction

LO Alpgen + PYTHIA

Total $\sigma$ normalized to Data
search for SUSY in acoplanar di-jet events
top quark: require 1 b tagged jet
Here there are staircases both with respect to the number of jets,

**DØ Run II Preliminary**

![Graph showing number of tagged events vs jet multiplicity]

- Data
- QCD
- W+light
- Wc
- Wcc
- Wbb
- ttbar → ll
- VV
- Z → ττ
- Single Top
- ttbar → jets

1 b-tag events
and with respect to the number of b-tagged jets.
CDF

Comparison of HT distributions between ttbar and W + jets events
ttbar/W+jets shape comparisons for 9 kinematic observables.
Using these and 10 more variables input to a neural network classifier, CDF has demonstrated the ability to observe $t\bar{t}$ events without $b$-tagging. Here are the last two steps in the staircase in that analysis.

\[
H_T = E_T + E_{T\ell} + \sum_i E_{Ti}
\]

**W+3 jets**

CDF Run2, 194 pb$^{-1}$

- $t\bar{t}$: $65.8 \pm 22.1^{+21.5}_{-21.5}$
- Multijet: $32.7 \pm 0.0^{+0.0}_{-0.0}$
- W-like: $419.6 \pm 28.9$

**W + 4 jets**

CDF Run2, 194 pb$^{-1}$

- $t\bar{t}$: $57.2 \pm 15.6^{+15.7}_{-15.7}$
- Multijet: $7.4 \pm 0.0^{+0.0}_{-0.0}$
- W-like: $52.8 \pm 16.8^{+14.9}_{-14.9}$
The CDF and D0 experiments reached an important milestone this past year with the observation of single top production.

This process has a rate about 10% of the rate for top quark pair production. It is actually two distinct processes, one with an s-channel pole, one with a t-channel pole:

The signature of single-top production is intermediate between those of $W + \text{jets}$ and top quark pair production.
The analyses are based on

\[ \ell + E_T + (2, 3) \text{ jet} \]

events with 1 b-tag.

To extract the single-top events from within these backgrounds, the CDF and DO events use automatic classifiers.

One method is to assign a weight to each event based on the lowest order matrix elements for the signal and background processes:

\[
D(x) = \frac{P(x|\text{signal})}{P(x|\text{signal}) + P(x|\text{background})}
\]

Other analyses use neural networks or boosted decision trees trained with Monte Carlo signal and background state.

I will show some figures from the matrix element based analyses.
First, apply the classifier to non-b tagged events.

2-jet

3-jet

s-channel
t-channel
then to the possible signal events:

- **s-channel**

- **t-channel**

2-jet events
3-jet events

s-channel

D0 Run II Preliminary

L = 0.9 fb⁻¹

Events

0 0.1 0.2 0.3 0.4 0.5 0.6 0.7 0.8 0.9 1

s-Disc Combined 3Jet

QCD
W + light jets
Wcc + jets
Wbb + jets
t\bar{t} \rightarrow lept + jets
t\bar{t} \rightarrow dilepton
signal: tb
signal: tqb
DATA

D0 Run II Preliminary

L = 0.9 fb⁻¹

Events

0 0.1 0.2 0.3 0.4 0.5 0.6 0.7 0.8 0.9 1

t-Disc Combined 3Jet

D0
same analysis, no b-tag
We can also move in the direction of multilepton signatures. Here there is another staircase, the **Baer-Tata staircase**.

Many new physics models such as supersymmetry predict 2, 3, 4 - lepton events in a steadily decreasing progression.

The Standard Model also produces such events, from **multiple heavy-quark decays** and jets faking leptons.

Fortunately, these come from the same W, Z, \(t\bar{t}\) + jets processes that we have already been discussing.

Electroweak backgrounds, e.g. 

\[ pp \rightarrow W^+ W^+ \rightarrow \ell^+ \ell^- + \text{jets} \]
signal cross sections from one of the models of Baer, Chen, Paige, Tata

$m_{\tilde{g}} \sim m_{\tilde{q}} \sim 750$ GeV

SS

Z + leptons

n(leptons)
Tevatron Run II $pp$ at $\sqrt{s} = 1.96$ TeV

thanks to M. Neubauer
These studies at the Tevatron give us confidence that we will be able to sort new physics events with heavy particles from the background due to Standard Model heavy particle production.

There is one genuinely new issue at the LHC. At the Tevatron, top quark pair production is at the pb level -- though already it is an important background in new physics searches. At the LHC, top quark pair production is at the nb level!

We need a way to obtain a relatively pure sample of $t\bar{t} + \text{jets}$ events to validate the theoretical models of this process used to estimate these backgrounds.
For example, in Padhi’s simulation of the single-lepton + MET signature:
One lepton Mode

W, t SM region

SUSY signal region

Padhi
Once we have convincingly established that new particles are produced at the LHC, it is another challenge to work out the properties of these particles and measure their masses and couplings.

This could fill another lecture, one that I will not have time for here. This is Prof. Nojiri’s subject, so I hope you can learn about it at another time.

In short, there are methods to find the masses of new particles at the 10% level, and some qualitative indications of spin and chirality assignments.
Instead of discussing this, I will rush into the future.

Tomorrow, we will talk about precision new particle spectroscopy at the ILC.